

# **$J^\pi$ from $(n,\gamma)$ Measurements and Statistical Model Calculations**

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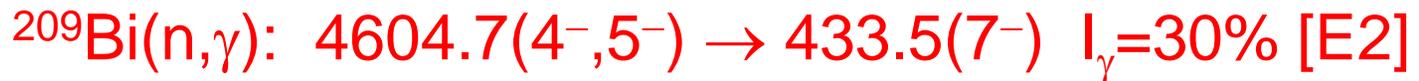
**Lawrence Berkeley National Laboratory**

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# $(n,\gamma)$ E=Thermal Data Considerations

- Thermal  $(n,\gamma)$  data are usually incomplete except for light ( $Z < 20$ ) isotopes. Beware of papers that normalize  $\gamma$ -ray intensities to  $\Sigma I_\gamma(\text{GS}) = 100$ .
- Primary  $\gamma$ -rays are mainly E1 or M1. E2  $\gamma$ -rays are much weaker

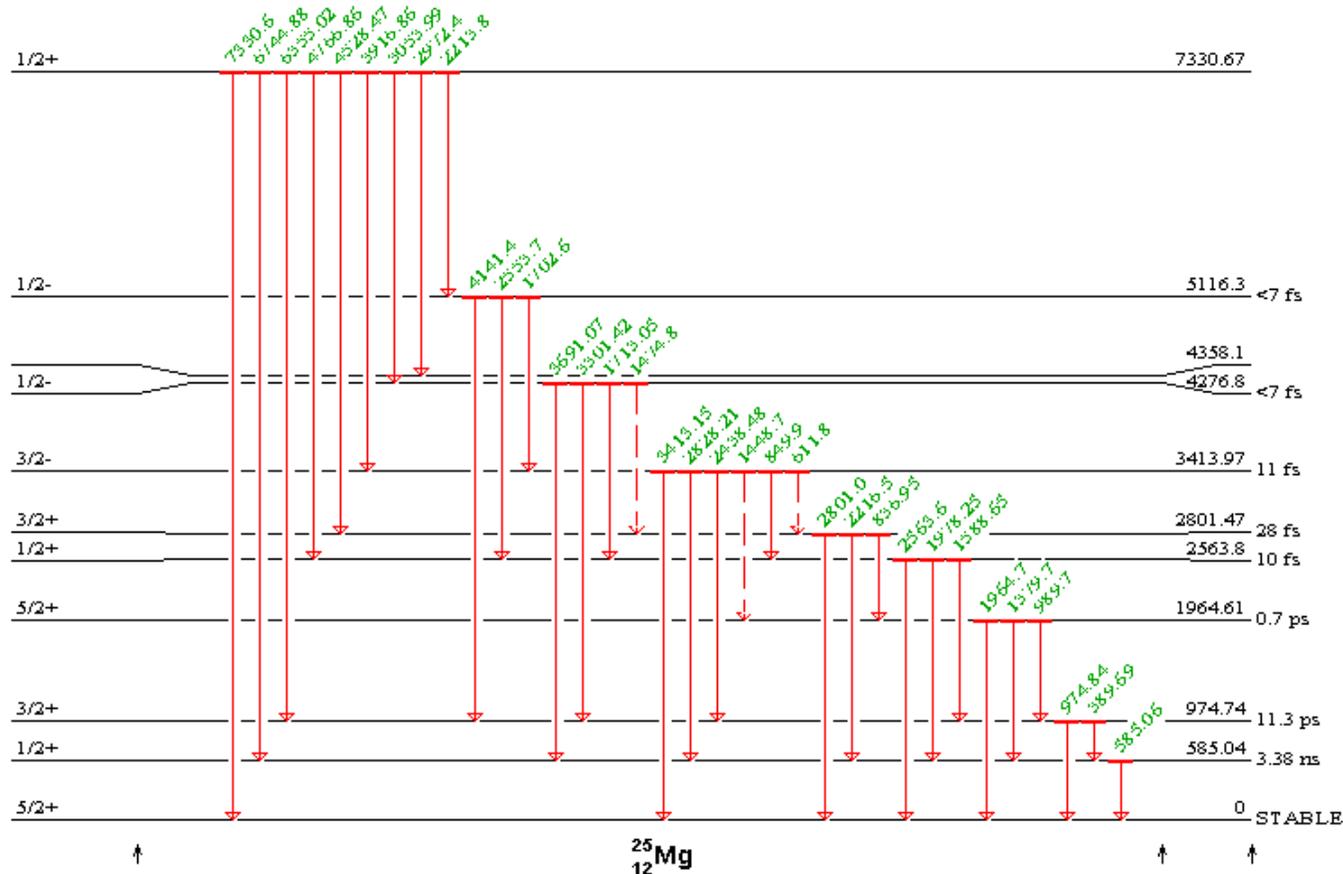


It is improbable that the strongest primary  $\gamma$ -ray is E2. Most likely  $J^\pi(433.5) = 6^-$  and the  $J^\pi$  assignments of low-lying levels in  $^{210}\text{Bi}$  are incorrect.

- Statistical model calculations can provide important information about spin assignments.

# Low-Z Data

The decay schemes of low-Z isotopes are often complete.



# Cross section balance for $^{24}\text{Mg}(n,\gamma)^{25}\text{Mg}$

Cross section balance for the  $^{25}\text{Mg}$  neutron capture decay scheme

E(Level)	$\sigma(\text{in})$	$\sigma(\text{out})$	$\Delta\sigma$
0	0.0536(14)	0.0	0
585.01(3)	0.0406(11)	0.0398(14)	0.0008(18)
974.68(3)	0.0157(4)	0.0158(4)	0.0001(6)
1964.69(10)	0.00022(2)	0.00026(3)	0.00004(4)
2563.35(4)	0.00202(10)	0.00179(7)	0.00023(12)
2801.54(9)	0.00047(4)	0.00061(5)	0.00013(6)
3413.35(3)	0.0411(14)	0.0416(11)	0.0005(18)
4276.33(4)	0.0105(4)	0.0107(3)	0.0002(5)
4358.2(5)	0.00009(2)	0.0	0.00009(2)
5116.37(15)	0.00038(4)	0.00027(3)	0.00011(5)
7330.53(4)	0.0	0.0539(14)	0.0539(14)
	$\sigma(\text{Mughabghab}[23])$	0.0536(15) b	
	$\sigma(\text{Measured, average})$	0.0538(14) b	

# High-Z Data

## $^{105}\text{Pd}(n,\gamma)^{106}\text{Pd}$ Level Feedings

E(level)	J <sup>π</sup>	Σσ <sub>γ</sub> (in)	Σσ <sub>γ</sub> (out)	ΔΣσ
0	0+	20.26		
511.844	2+	13.88	17.91	4.03
1128.04	2+	2.371	4.263	1.892
1133.79	0+	0.227	0.565	0.338
1229.2	4+	1.630	3.479	1.849
1557.67	3+	1.183	2.142	0.959
1562.16	2+	0.312	1.869	1.557
1706.44	0+	0.012	0.193	0.181
1909.39	2+	0.063	0.724	0.661
1932.37	4+	0.217	0.590	0.373
2001.56	0+	0.029	0.118	0.089
2077.1	6+	0.001	0.103	0.102
2077.37	(4)+	0.057	0.440	0.383
2084.39	-3	0.123	1.033	0.910
2242.4	2+	0.026	0.499	0.473
2278.47	0+	0	0.056	0.056
2282.89	4+	0.0007	0.275	0.274
2306.01	-3	0.053	0.542	0.489
2308.73	2+	0.000	0.283	0.283
2350.96	4+	0.018	0.304	0.286
2366.09	5+	0.003	0.116	0.114
2397.37	(5)-	0.055	0.263	0.209
2401	(2-,3-)	0.037	0.300	0.263
2439.11	2+	0.065	0.293	0.227
2472.09	0+	0.000	0.055	0.055
2484.76	(1-)	0.043	0.253	0.211
2500.01	-2	0.028	0.296	0.267
2578.64	(4-)	0.00004	0.221	0.221
...	...	...	...	...
...	...	...	...	...
9561.4	2+,3+		0.554	

For  $Z \geq 20$  measured neutron capture  $\gamma$ -ray decay schemes are generally incomplete due to unresolved continuum  $\gamma$ -rays.

$$\sigma_0 = 21.0 \pm 1.5 \text{ b (Mughabghab)}$$

$$\sigma_\gamma(\text{primary } \gamma\text{-rays}) = 0.55 \text{ b}$$

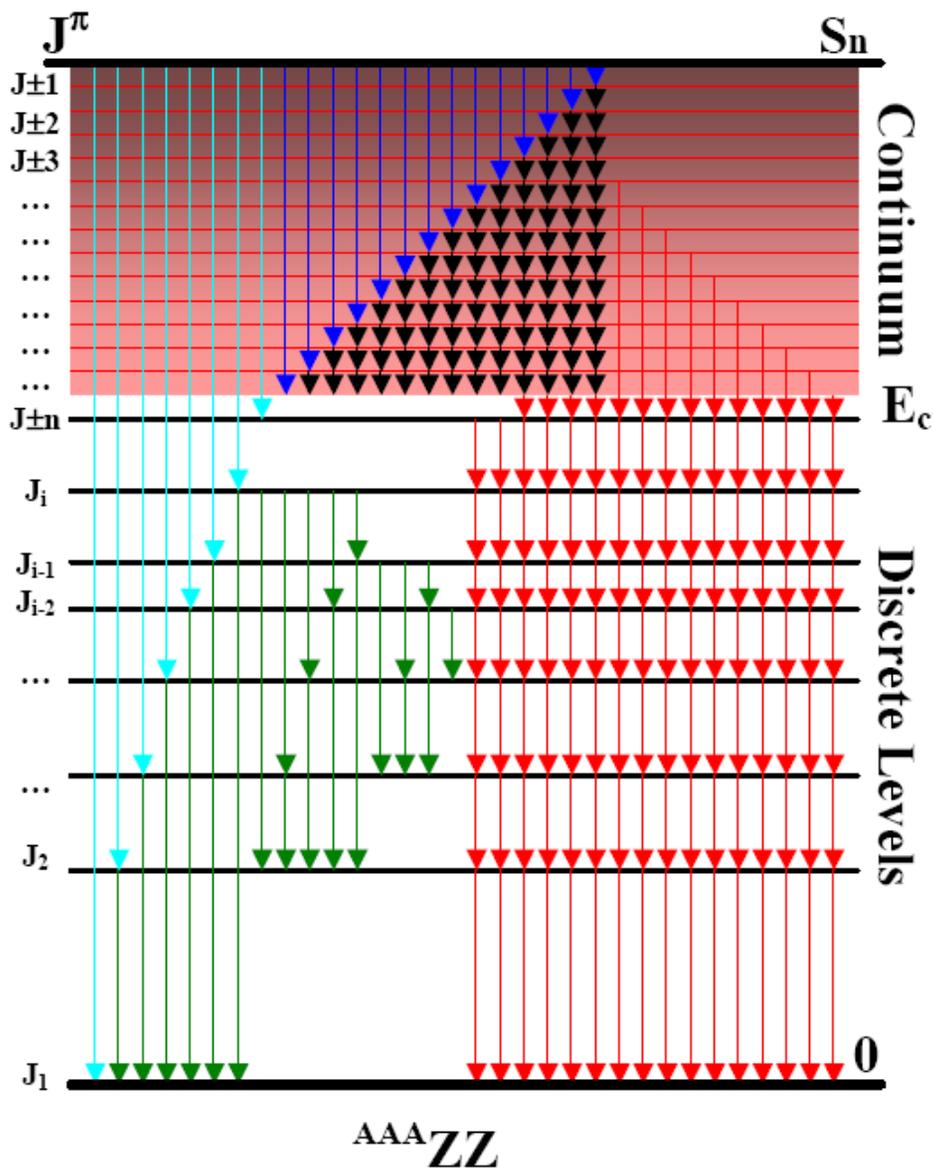
$$\sigma_\gamma(\text{secondary } \gamma\text{-rays, observed}) = 20.3 \pm 0.3 \text{ b}$$

$$\underline{\sigma_\gamma(\text{secondary } \gamma\text{-rays, statistical}) = 1.4 \pm 0.3 \text{ b}}$$

$$\sigma_0 = 21.7 \pm 4 \text{ b}$$

$^{105}\text{Pd}(n,\gamma)^{106}\text{Pd}$  level feeding cross sections calculated from EGAF data.

# Statistical Model Calculations



The  $(n, \gamma)$  continuum feeding is statistical so it can be calculated if

1.  $\sigma_\gamma$  deexciting levels below a cutoff energy  $E_{crit}$  is complete. ↓
2. Primary  $\sigma_\gamma$  populating the levels below  $E_{crit}$  from the capture state is complete. ↓
3.  $J^\pi$  of levels below  $E_{crit}$  are well known.
4. Level density  $\rho(E > E_{crit}, J)$  is known.
5. Photon strength  $f(E_\gamma)$  deexciting levels above  $E_{crit}$  is known ↓ ↓

# DICEBOX Calculations

DICEBOX is a Monte Carlo code by F. Becvar and M. Krlicka. It generates complete simulated neutron capture decay schemes constrained by known nuclear properties and statistical models.

- A. Discrete primary and secondary  $\sigma_\gamma$  data from EGAF
- B.  $J^\pi$  data for  $E < E_{\text{crit}}$  from Reaction Input Parameter Library (RIPL)

C. Level density models

- 1. Constant temperature (CT)
- 2. Back-shifted Fermi (BSF) model

$$\rho(E, J) = \frac{f(J)}{T} \exp\left(\frac{E - E_0}{T}\right)$$

$$\rho(E, J) = f(J) \frac{\exp\left(2\sqrt{a(E - E_1)}\right)}{12\sqrt{2}\sigma_c a^{1/4}(E - E_1)^{5/4}}$$

D. EI Photon Strength

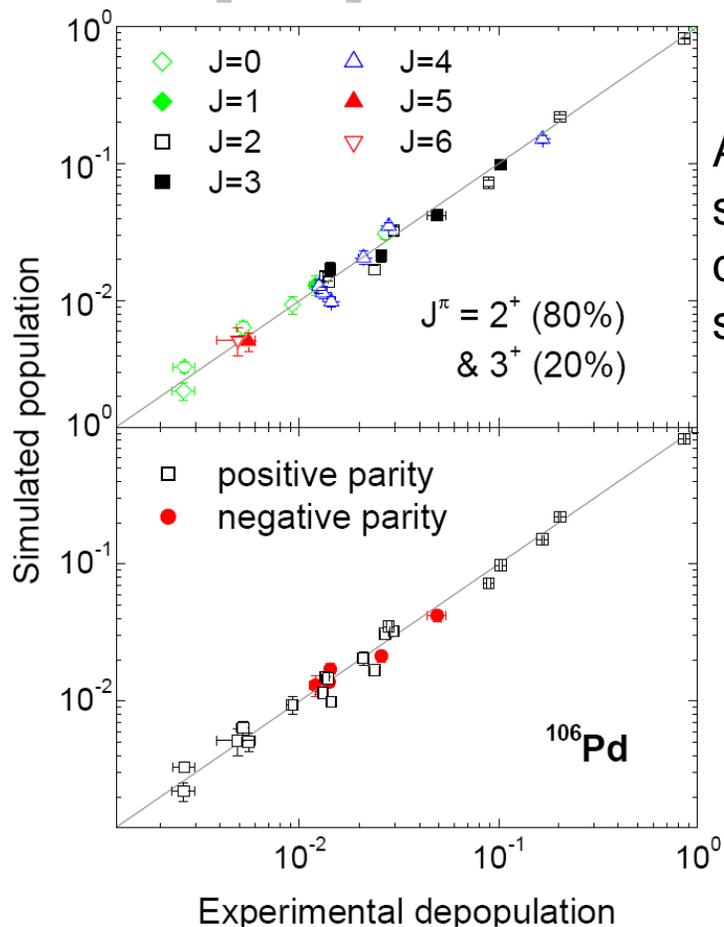
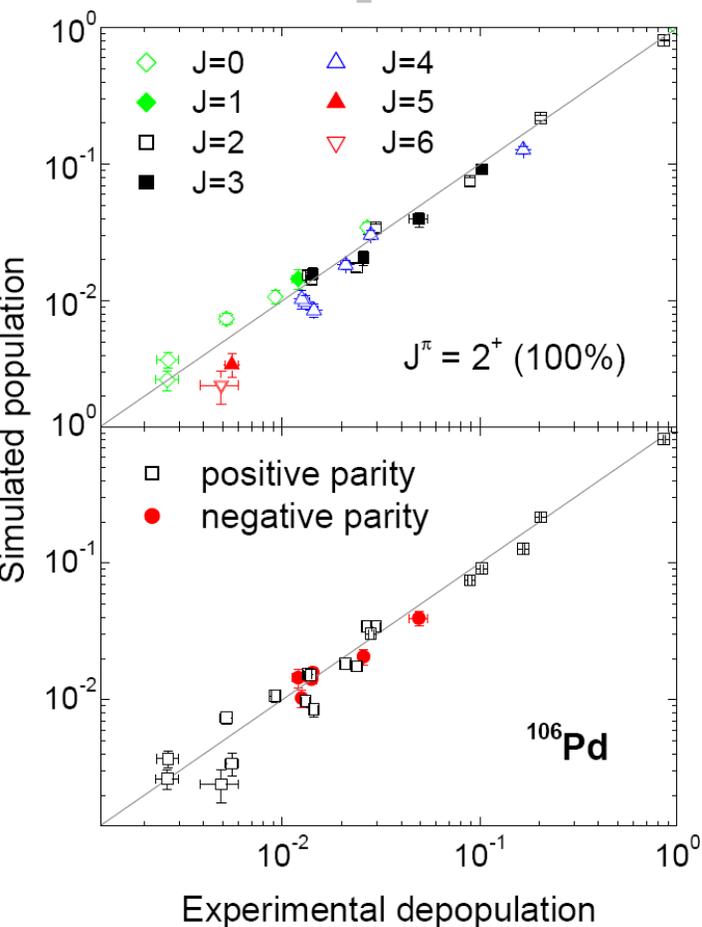
- 1. Brink-Axel (BA)
- 2. Kadmenski, Markushev, Furman (KMF) for spherical nuclei
- 3. Kopecky *et al* generalized Laurentian (GLO), temperature dep.

$$f_{\text{BA}}^{(E1)}(E_\gamma) = \frac{1}{3(\pi\hbar c)^2} \frac{\sigma_G E_\gamma \Gamma_G^2}{(E_\gamma^2 - E_G^2)^2 + E_\gamma^2 \Gamma_G^2}$$

E. M1 Photon Strength

- 1. Single Particle (SP),  $f^{(E1)}/f^{(M1)}=5-7$  or  $f(M1)=1.2 \times 10^{-8} \text{ MeV}^{-3}$
- 2. Spin-Flip (SF), *Laurentzian resonance*  $\approx 8.5 \text{ MeV}$ ,  $\Gamma_{\text{SF}} \approx 4 \text{ MeV}$

# Population/Depopulation Plot

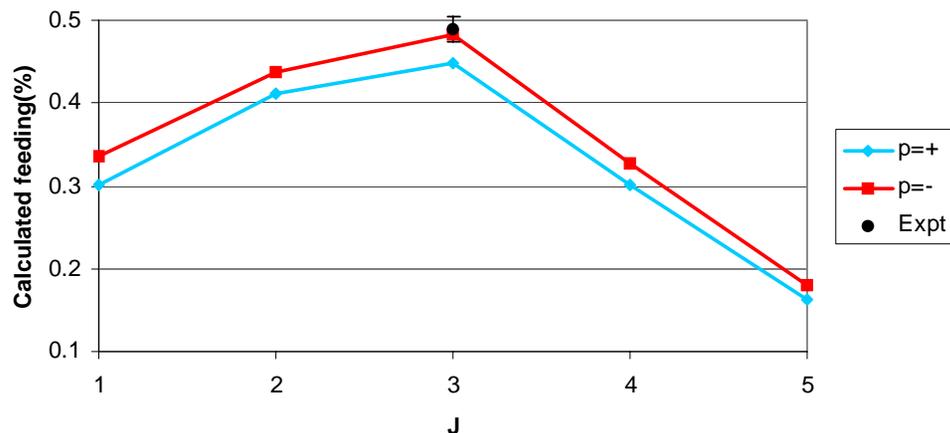


Admixture of capture state  $J^\pi$  values determined by least-squares fit.

Significant deviations from the line can be indications of problems with input RIPL nuclear structure data. Either  $J^\pi$  or  $\gamma$ -ray branching may be wrong.

# Data Problems: $J^\pi$ Errors

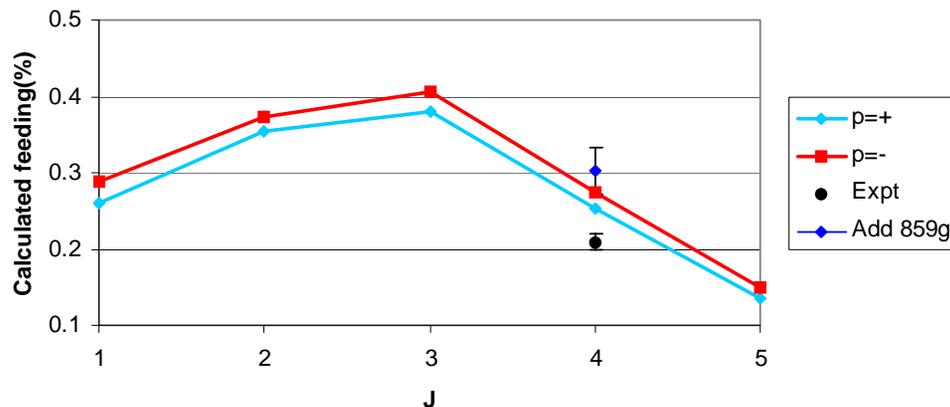
Statistical Feeding to 2306 Level in  $^{106}\text{Pd}$  (4- in ENSDF)



Statistical model calculations can be used to constrain  $J^\pi$  values.

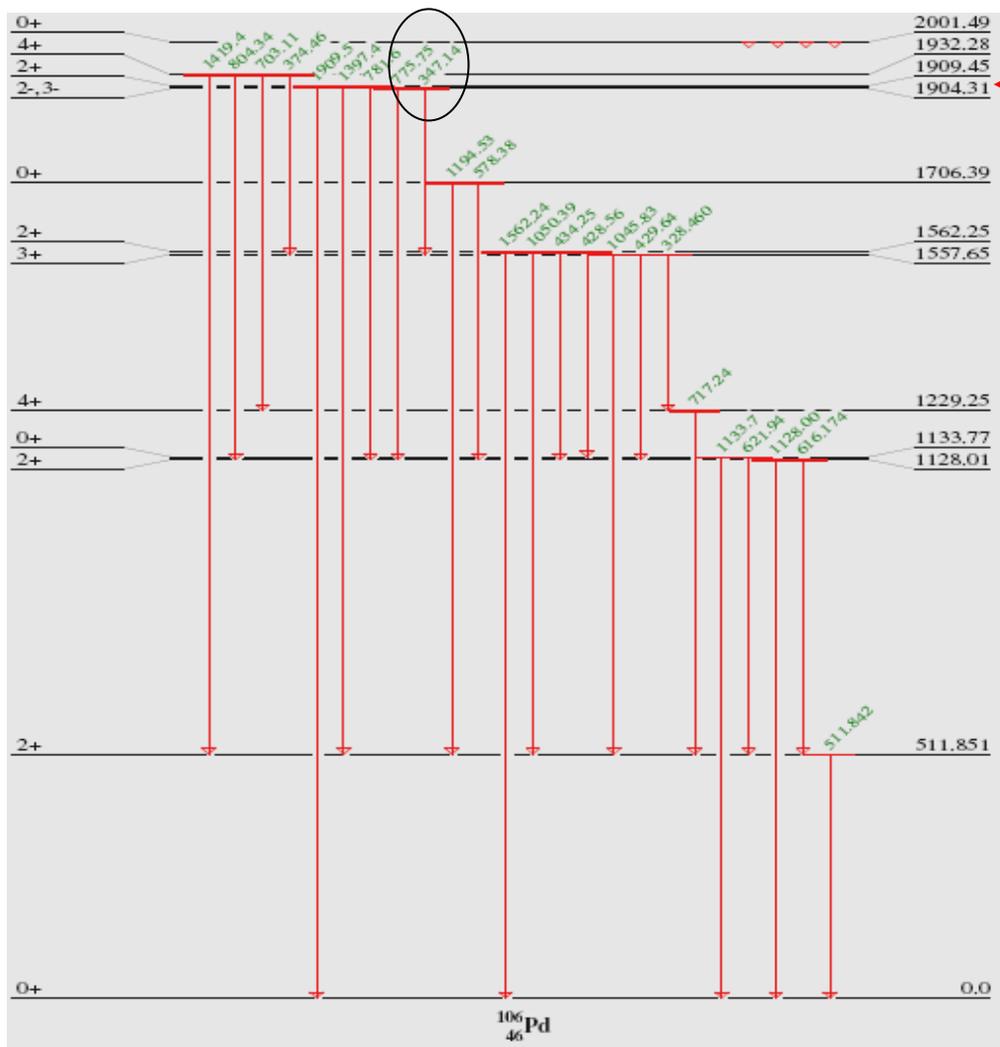
2306-keV level feeding is consistent with  $J^\pi=3^-$  not  $J^\pi=4^-$  adopted in ENSDF on basis of  $\gamma\gamma(\theta)$ .

Statistical Feeding to 2397 Level in  $^{106}\text{Pd}$  (5- in ENSDF)



2397-keV level feeding is consistent with  $J^\pi=4^-$  when additional  $\gamma$ -ray is placed. Assignment of  $J^\pi=5^-$  in ENSDF is based on  $L=(5)$  in  $(p,t)$ .

# Level Scheme Errors



**Mistaken level assignment**

1904.3 keV level assigned using the Ritz principal.

$$\sigma_{\gamma}(1904)_{\text{expt}} = 0.12 \text{ b}$$

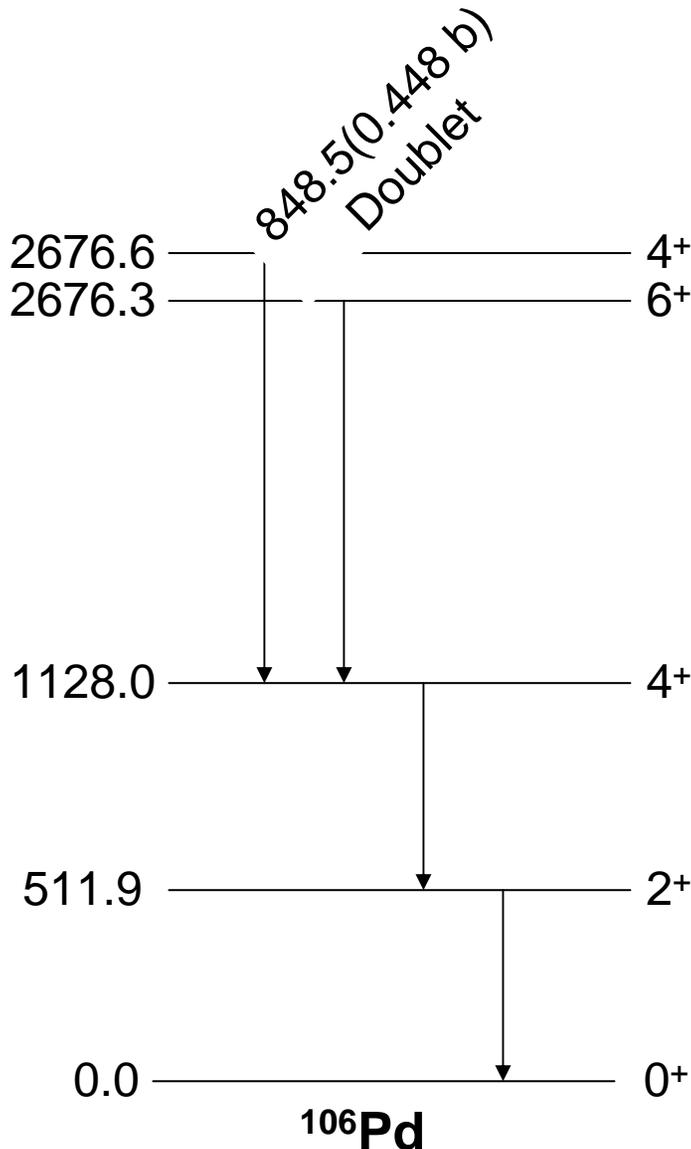
$$\sigma_{\gamma}(1904)_{\text{DICEBOX}} = 1.13 \text{ b}$$

Reassigning placement of the 347- and 776-keV  $\gamma$ -rays to deexcite the 1909.4-keV level gives

$$\sigma_{\gamma}(1909)_{\text{expt}} = 0.62 \text{ b}$$

$$\sigma_{\gamma}(1909)_{\text{DICEBOX}} = 0.83 \text{ b}$$

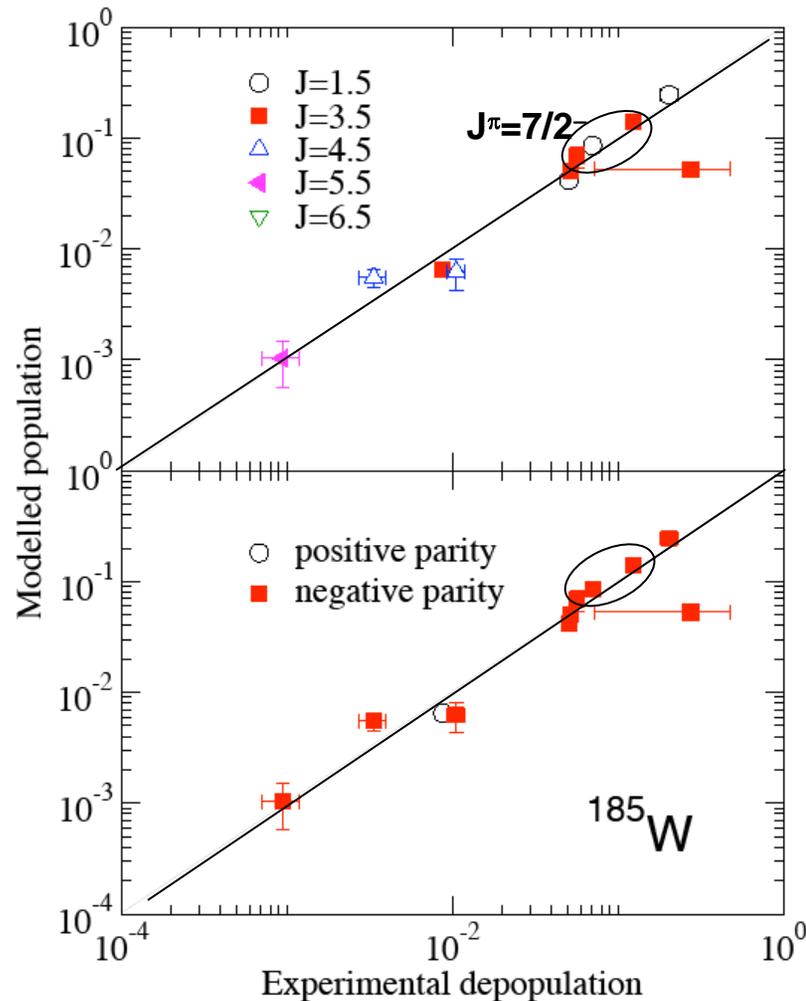
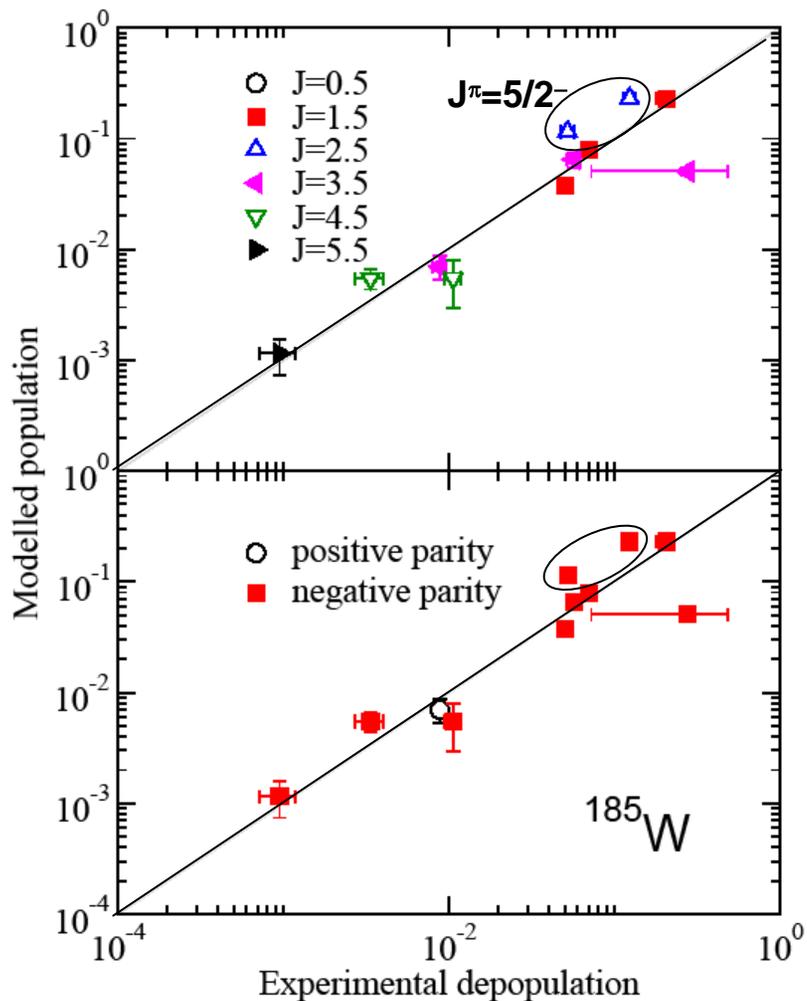
# Resolving $\gamma$ -ray Doublets



Doublet cross section ( $\sigma_{848.5\gamma} = 0.448 \text{ b}$ ) populating the  $4^+$  and  $6^+$  can be divided based on DICEBOX calculation.

<b>E<sub>g</sub></b> <b>(keV)</b>	<b>Transition</b>	<b>Calculated</b> <b>cross section</b> <b>(barns)</b>
848.6	$4^+ \rightarrow 4^+$	$0.345 \pm 0.022$
848.3	$6^+ \rightarrow 4^+$	$0.103 \pm 0.022$

DICEBOX Population/Depopulation plot for  $^{184}\text{W}(n,\gamma)^{185}\text{W}$  data. Levels at 66- and 188-keV have  $J^\pi=5/2^-$  in ENSDF. Statistical model calculations indicate that both levels are more consistent with  $7/2^-$ . Many other “definite”  $J^\pi$  values in  $^{185}\text{W}$  must be revised.



# $^{185}\text{W}$ $J^\pi$ Justifications

Level	$J^\pi$ (ENSDF)	ENSDF $J^\pi$ justification	New $J^\pi$
0.0	$3/2^-$	Atomic beam	$3/2^-$
23.5	$1/2^-$	$1/2^-, 3/2^-$ (n, $\gamma$ ), DWBA	$3/2^-$
65.9	$5/2^-$	L=3 in (d,t),(d,p), M1+E2 to $3/2^-$ $\alpha(L)_{\text{exp}}=10(2), 11(3)$ . $\alpha(L)_{M1}=2.11$ , $\alpha(L)_{E2}=16.6$ $\alpha(M)_{\text{exp}}=3.5(15)$ . $\alpha(M)_{M1}=0.48$ , $\alpha(M)_{E2}=4.2$ ICC(65.9 $\gamma$ ) is also consistent with E2	$7/2^-$
93.3	$3/2^-$	L=1 in (d,t),(d,p)	$3/2^-$
173.7	$7/2^-$	M1+E2 to $5/2^-$ , E2 to $3/2^-$ (GS)	$7/2^-$
187.9	$5/2^-$	L=3 in (d,t),(d,p), $\gamma$ -ray to $1/2^-$	$7/2^-$

$J=L+1/2$  appears to be favored in (d,t),(d,p) reactions

# Conclusions

- Statistical model calculations provide an important check of  $J^\pi$  assignments
- “Definite”  $J^\pi$  assignments may be wrong
- New publications using statistical  $J^\pi$  arguments are coming
- Simple statistical arguments should be used by evaluators when adopting  $J^\pi$  values.
- An evaluator code for doing statistical calculations would be useful.