



The Short-Collision-Time Approximation in ENDF

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Summary I

- The short-collision-time approximation (SCTA) is used in ENDF to extend the range of tabulated values of the incoherent inelastic cross section to higher values of α and β than can be calculated from the tabulated thermal neutron scattering law $S(\alpha, \beta, T)$.
- It's usually attributed to Lamb (PR 55, 190, 1939) and paraphrased “at high values of energy and momentum transfer, the neutron scatters from a bound-atom as it would from a free gas at an effective temperature.”
- However, the formula introduced by the General Atomics team in the late sixties when they developed the ENDF files differs from that description, a fact which has led to considerable “discussion”.



Summary II

- We will display the GA formula, and for the record, five “alternative” forms which appear in ENDF-related documents:
 - “The NJOY Nuclear Data Processing System”, LA-9303-M, Vol II (ENDF-324), May 1982 (currently available from RSIC)
 - “New Thermal Neutron Scattering Files for ENDF/B-VI Release 2”, LA-12639-MS (ENDF 356), March 1994 (currently available from LANL)
 - “ENDF- 6 Formats Manual, Data Formats and Procedures for the Evaluated Nuclear Data File ENDF/B-VI and ENDF/B-VII”, ENDF-102, BNL-NCS-44945-05-Rev June 2005 (currently available from NNDC)
 - Two “corrected” versions from the Formats Committee which have not yet been issued.



Introduction I

- The double-differential cross section in terms of the symmetrized Scattering Law, $S(\alpha, \beta, T)$, is (in ENDF notation)

$$d^2\sigma(E \rightarrow E', \mu, T)/dE'd\Omega = (\sigma_b/4\pi T) \sqrt{(E'/E)} \exp(-\beta/2) S(\alpha, \beta, T)$$

in barns/steradian-eV with $\alpha = [E + E' - 2\mu\sqrt{EE'}]/AT$ and $\beta = (E' - E)/T$ the dimensionless momentum-squared and energy transfers. T is the temperature in eV and $A = M/m$.

- Detailed balance “DB” is defined (for our purposes) by:

$$M(E, T)\sigma(E \rightarrow E') = M(E', T)\sigma(E' \rightarrow E)$$

where

$$M(E, T) = (E/T^2) \exp(-E/T)$$



Introduction II

- DB is necessary and sufficient for the neutrons to achieve thermal equilibrium with the moderator. It is required of all ordinary moderators.
- Symmetry of $S(\alpha, \beta, T)$ in beta is necessary and sufficient for the **cross section** to obey DB. That is what the factor $\exp(-\beta/2)$ accomplishes on the preceding slide.
- DB is a thermodynamic statement about media in thermodynamic equilibrium. As such, it involves the ambient, thermodynamic temperature, regardless of what pseudo-temperatures may occur inside of $S(\alpha, \beta, T)$.



Introduction III

- The monatomic ideal gas Scattering Law is

$$S_{MIG}(\alpha, \beta, T) = [1/\sqrt{4\pi\alpha}] \exp\{-(\alpha^2 + \beta^2)/4\alpha\}$$

- It is symmetric in β , so the cross section satisfies DB at the temperature T . The cross section is

$$d^2\sigma_{MIG}(E \rightarrow E', \mu, T)/dE'd\Omega = (\sigma_b/4\pi T) \sqrt{(E'/4\pi\alpha E)} \exp\{-(\alpha + \beta)^2/4\alpha\}$$

- If it is integrated over angles, the “total” cross section is

$$\sigma_{MIG}(E \rightarrow E') = (\sigma_b/E) \operatorname{erf}(\sqrt{E'/T}) \quad E' < E \quad \text{downscatter}$$

$$\sigma_{MIG}(E \rightarrow E') = (\sigma_b/E) \operatorname{erf}(\sqrt{E/T}) \exp\{-(E' - E)/T\} \quad E' > E \quad \text{upscatter}$$



Introduction IV

- Since we are interested in high E and E'-values, the erf's are flat and the interesting thing is the upscattering tail $\exp\{-E'/T\}$ that “drops off” more slowly as the temperature increases.

- For example, at a higher T*, the MIG cross section would be

$$d^2\sigma(E \rightarrow E', \mu, T^*)/dE'd\Omega = (\sigma_b/4\pi T^*) \sqrt{E'/4\pi\alpha^*E} \exp\{-\frac{(\alpha^* + \beta^*)^2}{4\alpha^*}\}$$

and would upscatter more, $\exp\{-E'/T^*\}$



The GA Formula I

- The reference for this formula is “Neutron Scattering Kernels Calculations at Epithermal Energies”, GA-9950, UC-80, Reactor Technology, March 17, 1970, G. M. Borgonovi, Gulf General Atomic, Inc.
- In that report, equation (1) defines the double-differential cross section, and is the same definition as the one shown previously. It has a typographical error, in that the cross section is differential in dE although there is no E in the formula.
- Equation (2) has a typographical error. The symbol κ is used where β should be.
- Equation (3) has a typographical error. The symbol $S(\kappa, \omega)$ is used where “S-bar” should be.



The GA Formula II

- To derive the SCTA, the following assumptions were made:
 - incoherent scattering
 - Gaussian approximation for the intermediate scattering function
 - Cubic Bravais crystal, harmonic motion
- The SCTA itself is “ $\sin \omega t = \omega t$; $\cos \omega t = 1 - (\omega t)^2 / 2$ ”
- The result is that the Scattering Law can be written as their Equation (13):

$$S_{sc}(\alpha, \beta, T) = [1 / \sqrt{(4\pi\alpha T^* / T)}] \exp\{-(\alpha + \beta)^2 T / 4\alpha T^* - \beta / 2\},$$

- Although this does not have the free-gas form shown on the preceding slide, when substituted into the cross section formula, it produces a free-gas-at- T^* , which was shown on Slide 3.



The GA Formula III

- That has two problems:
 - Since it satisfies DB at T^* , it doesn't satisfy it at T , where we want it.
 - It has an unphysically long upscatter tail, as shown on the following slide.
- Both of these defects can be remedied by re-defining the Scattering Law, which they did as follows:

$$S_{sc}(\alpha, \beta, T) ==>$$

$$S_{sc}(\alpha, -|\beta|, T) = [1/\sqrt{(4\pi\alpha T^*/T)}] \exp\{-[(\alpha + |\beta|)^2 T/4\alpha T^* - |\beta|/2]\}$$

- This is the GA formula for the SCTA. It is currently coded that way in THERMR. It is written that way in only one document that I have seen, namely "Thermal Neutron Scattering Data, etc.", INDC(NDS)-0470, M. Mattes and J Keinert, April 2005. However, Andrej and Mike assure me that it is OK in the unpublished forthcoming revision to ENDF-102.

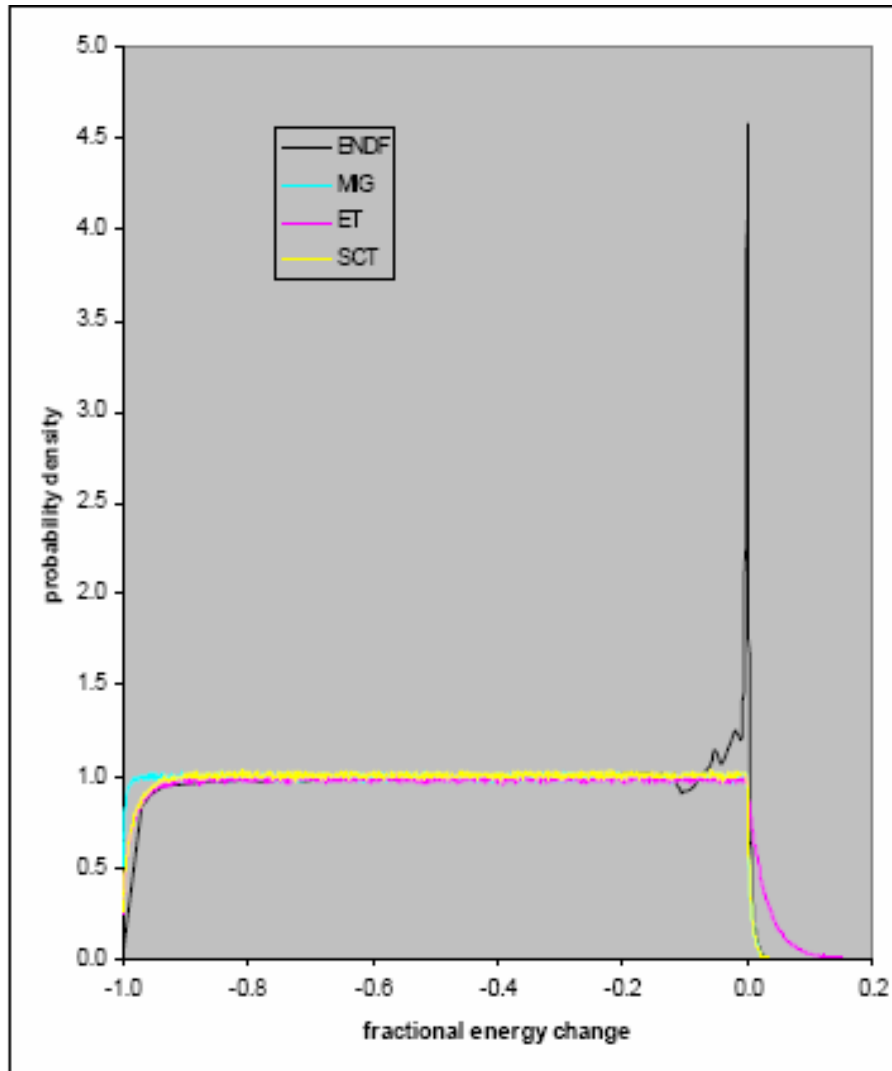


Figure 1. Fractional Energy Change PDF for 4.09 eV Neutrons on ^1H in Water at 293.6 K

- ENDF is LEAPR + GA SCTA.
- MIG is FREE GAS at T-ambient
- ET is FREE GAS at T-effective
- SCT is GA SCTA



Some “Alternative” Forms I

- The reason for displaying these is to provide a definitive statement of what GA intended, and to identify known misprinted versions unambiguously.
- The square root in front is usually OK so here are just the exponents without the minus sign:

1. “The NJOY Nuclear Data Processing System”, LA-9303-M, Vol II (ENDF324), May 1982 (currently available from RSIC)

$$(|\alpha - \beta)^2 T / 4 |\alpha| T^* + |\beta| / 2$$

Missing absolute-value signs on first β ; unnecessary absolute-value signs on α .



Some “Alternative” Forms II

2. “New Thermal Neutron Scattering Files for ENDF/B-VI Release 2”, LA-12639-MS (ENDF 356), March 1994 (currently available from LANL)

$$(\alpha - |\beta|)^2 T / \alpha T^* + |\beta|/2$$

Missing factor of 4 in denominator.

3. “ENDF- 6 Formats Manual, Data Formats and Procedures for the Evaluated Nuclear Data File ENDF/B-VI and ENDF/B-VII”, ENDF-102, BNL-NCS-44945-05-Rev June 2005 (currently available from NNDC)

$$(\alpha - \beta)^2 T / 4\alpha T^* + |\beta|/2$$

Missing absolute-value signs on first β



Some “Alternative” Forms III

4. A “corrected” version from the (former) Formats Committee which has not yet been issued.

$$(\alpha - \beta)^2 T / 4\alpha T^* + (\beta + |\beta|) / 2$$

Missing absvals on the first β ; an extra β probably from the cross section.

5. A “corrected” version of the above, which has not yet been issued.

$$(\alpha - |\beta|)^2 T / 4\alpha T^* + (\beta + |\beta|) / 2$$

An extra β probably from the cross section.



Acknowledgement

Material on the GA formula was taken from a paper submitted to the International Conference on Mathematics, Computational Methods & Reactor Physics (M&C 2009) Saratoga Springs, NY, May 3 – 7, 2009, entitled “Comparison of Some Monte Carlo Models for Bound Hydrogen Scattering”, by TM Sutton, TH Trumbull, and CR Lubitz.