## ${ }^{93}$ Rh $\varepsilon$ decay 2004De40

| Type | Author | History <br> Full Evaluation$\frac{\text { Citation }}{\text { Coral M. Baglin }}$ |  |
| :---: | :---: | :---: | :---: |

Parent: ${ }^{93} \mathrm{Rh}: \mathrm{E}=0.0 ; \mathrm{J}^{\pi}=\left(9 / 2^{+}\right) ; \mathrm{T}_{1 / 2}=11.9 \mathrm{~s} 7 ; \mathrm{Q}(\varepsilon)=8203.06 ; \% \varepsilon+\% \beta^{+}$decay $=100.0$
${ }^{93} \mathrm{Rh}-\mathrm{Q}(\varepsilon)$ : From 2009AuZZ; Q=8090 410 from systematics (2003Au03).
${ }^{93} \mathrm{Rh}-\mathrm{T}_{1 / 2}$ : Calculated from a fit to the sum of the individual time-to-digital spectra gated on each of the 7 transitions attributed to
${ }^{93}$ Rh decay. Half-life was accurately measured using a macrocycle of a beam-on period followed by a beam-off period. The on/off times were altered to suit the expected $\mathrm{T}_{1 / 2}$ of the isotope of interest. The tdc was started at the beginning of each macrocycle, recording the time of each triggered event relative to the start.
2004De40: ${ }^{93} \mathrm{Rh}$ source from ${ }^{58} \mathrm{Ni}\left({ }^{40} \mathrm{Ar}^{11+}, \mathrm{P} 4 \mathrm{~N}\right), \mathrm{E}=171 \mathrm{MeV}$ At target face (after degradation of 250 MeV beam using Ta foils); $99.93 \%{ }^{58} \mathrm{Ni}$ target; recoils were stopped and neutralized in 500 mbar of purified Ar gas before being ionized selectively (according to Z ) using two dye lasers tuned to the resonant atomic transitions of Rh to enhance ionization, and thus extraction; laser-ionized nuclei guided towards the LISOL mass separator by a sextupole ion guide, then implanted onto movable tape; $\beta$-sensitive plastic $\Delta \mathrm{E}$ detectors; 2 HPGe detectors; measured $\mathrm{E} \gamma, \mathrm{I} \gamma, \gamma \gamma$ coin, $\beta \gamma$ coin, $\mathrm{I} \beta$, isotope $\mathrm{T}_{1 / 2}$.
No evidence (neither the IT nor any $\beta$-delayed $\gamma$ events) was found by 2004De40 for the presence of the known (1976De37) $\mathrm{T}_{1 / 2}=10.8 \mathrm{~s}, 1 / 2^{-}$isomer In ${ }^{93} \mathrm{Ru}$.

## ${ }^{93}$ Ru Levels

2004De40 conclude that many of the low-lying ${ }^{93} \mathrm{Ru}$ states populated in ${ }^{93} \mathrm{Rh} \varepsilon$ decay can be understood as belonging to $\pi\left(\mathrm{p}_{1 / 2}, \mathrm{~g}_{9 / 2}\right)^{-6} \nu \mathrm{~g}_{9 / 2}^{-1}$ configurations.

| $\mathrm{E}\left(\right.$ level) ${ }^{\dagger}$ | $\mathrm{J}^{\text {¢ }}$ | Comments |
| :---: | :---: | :---: |
| 0.0 | $(9 / 2)^{+}$ |  |
| 1359.4210 |  | Additional information 1. |
| 1393.3120 | $(13 / 2)^{+}$ | Additional information 2. |
| 1629.9210 |  | Additional information 3. |
| 1842.13 |  | Additional information 4. |
| 2273.5314 |  | Additional information 5. |
|  |  | $\mathrm{J}^{\pi}: 2004 \mathrm{De} 40$ conclude that this level is a likely candidate for the first excited $9 / 2^{+}$state predicted at 2 MeV by shell-model calculations. |

${ }^{\dagger}$ From least-squares fit to E $\gamma$.
$\ddagger$ From Adopted Levels.

| $\underline{\varepsilon, \beta^{+} \text {radiations }}$ |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| E(decay) | $\underline{\text { E(level) }}$ | $\mathrm{I} \beta^{+}$\# | $1 \varepsilon^{\# \#}$ | $\underline{\log f_{t}{ }^{\ddagger}}$ | $\underline{\mathrm{I}\left(\varepsilon+\beta^{+}\right)^{\dagger \#}}$ | Comments |
| (5929.5 \%) | 2273.53 | 5.110 | 0.153 | 5.689 | 5.210 | $\begin{aligned} & \text { av } \mathrm{E} \beta=2266.43 \text { 30; } \varepsilon \mathrm{K}=0.024271 \text { 9; } \varepsilon \mathrm{L}=0.002962 \mathrm{l} \text {; } \\ & \varepsilon \mathrm{M}+=0.00069613 \end{aligned}$ |
| (6360.9 7) | 1842.1 | 2.99 | 0.06620 | 6.0914 | 3.09 | $\begin{aligned} & \text { av } \mathrm{E} \beta=2474.6733 ; \varepsilon \mathrm{K}=0.0190517 ; \varepsilon \mathrm{L}=0.00232419 \text {; } \\ & \varepsilon \mathrm{M}+=0.00054612 \end{aligned}$ |
| (6573.1 \%) | 1629.92 | 3.28 | 0.06516 | 6.1211 | 3.38 | $\begin{aligned} & \text { av } \mathrm{E} \beta=2577.3430 ; \varepsilon \mathrm{K}=0.0170266 ; \varepsilon \mathrm{L}=0.00207677 \text {; } \\ & \varepsilon \mathrm{M}+=0.00048802 \end{aligned}$ |
| (6809.7 ©) | 1393.31 | 4.18 | 0.07314 | 6.109 | 4.28 | $\text { av } \mathrm{E} \beta=2692.02 \text { 31; } \varepsilon \mathrm{K}=0.0150935 ; \varepsilon \mathrm{L}=0.0018406 \text { 6; }$ $\varepsilon \mathrm{M}+=0.00043252$ |
| (6843.6 6) | 1359.42 | 4.59 | 0.07915 | 6.079 | 4.69 | $\log f t$ : value is unrealistically low compared to that expected for a $\Delta \mathrm{J}=2, \Delta \pi=\mathrm{No} \varepsilon$ transition. <br> av $\mathrm{E} \beta=2708.46$ 30; $\varepsilon \mathrm{K}=0.0148415 ; \varepsilon \mathrm{L}=0.00180976$; $\varepsilon \mathrm{M}+=0.00042522$ |
| (8203.0 6) | 0.0 | 7916 | 0.7415 | 5.269 | 8016 | $\begin{aligned} & \text { av } \mathrm{E} \beta=3370.72 ; \varepsilon \mathrm{K}=0.0080752 ; \varepsilon \mathrm{L}=0.00098383 ; \\ & \varepsilon \mathrm{M}+=0.0002311 \end{aligned}$ |

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## ${ }^{93} \mathrm{Rh} \varepsilon$ decay 2004De40 (continued)

$$
\varepsilon, \beta^{+} \text {radiations (continued) }
$$

${ }^{\dagger}$ Absolute intensity determined by 2004De40 from $\mathrm{I}\left(\gamma^{ \pm}\right)$after correction for contribution from decay of the ${ }^{93} \mathrm{Ru}$ isobar. All $\mathrm{I}\left(\gamma^{ \pm}\right)$not associated with $\gamma$ events visible in the $\gamma$-ray spectra were attributed to the g.s. branch. Consequently, this branch may be overestimated because it will incorporate branching to states whose deexciting $\gamma$-rays are too weak or too energetic ( $>4 \mathrm{MeV}$ ) to have been detected in this experiment. Also, the existence of such transitions may result in an under-estimation of branching to some excited states.
\# Values should probably be regarded as lower limits because the large Q value suggests the possibility of significant unobserved feeding to highly excited states which subsequently decay to low-lying levels. Weakly populated states and branching to states producing $\gamma$-rays outside the $4 \mathrm{MeV} \gamma$-energy range may be wrongly attributed to ground-state decay. A further consequence may be that apparently-forbidden $\beta$ decays may result from $\gamma$ transitions from higher-lying levels fed by allowed $\varepsilon$ transitions.
\# Absolute intensity per 100 decays.

$$
\gamma\left({ }^{93} \mathrm{Ru}\right)
$$

I $\gamma$ normalization: from comparison of $\Sigma(\mathrm{I}(\gamma+\mathrm{ce})$ to $\mathrm{g} . \mathrm{s}$.$) with authors' summed \% \varepsilon+\% \beta^{+}$to excited states (=20.3 20).

| $\mathrm{E}_{\gamma}$ | $\mathrm{I}_{\gamma}{ }^{\dagger}$ | $\mathrm{E}_{i}$ (level) | $\mathrm{J}_{i}^{\pi}$ | $\mathrm{E}_{f}$ | $\mathrm{J}_{f}^{\pi}$ | Mult. ${ }^{\dagger}$ | $\alpha^{\#}$ | Comments |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 482.63 | 207 | 1842.1 |  | 1359.42 |  |  |  |  |
| 643.61 | 428 | 2273.53 |  | 1629.92 |  |  |  |  |
| 1359.41 | 100 | 1359.42 |  | 0.0 | (9/2) ${ }^{+}$ |  |  |  |
| 1393.32 | 7214 | 1393.31 | $(13 / 2)^{+}$ | 0.0 | (9/2) ${ }^{+}$ | E2 | 0.0004837 | $\begin{aligned} & \alpha=0.0004837 ; \alpha(\mathrm{K})=0.0003826 ; \\ & \quad \alpha(\mathrm{L})=4.33 \times 10^{-5} 6 ; \alpha(\mathrm{M})=7.93 \times 10^{-6} \quad 11 ; \\ & \alpha(\mathrm{N}+. .)=4.98 \times 10^{-5} 7 \\ & \alpha(\mathrm{~N})=1.283 \times 10^{-6} 18 ; \alpha(\mathrm{O})=6.81 \times 10^{-8} \quad 10 ; \\ & \alpha(\mathrm{IPF})=4.84 \times 10^{-5} 7 \end{aligned}$ |
| 1629.91 | 9819 | 1629.92 |  | 0.0 | (9/2) ${ }^{+}$ |  |  |  |
| 1842.46 | 3113 | 1842.1 |  | 0.0 | (9/2) ${ }^{+}$ |  |  |  |
| 2273.89 | 4813 | 2273.53 |  | 0.0 | (9/2) ${ }^{+}$ |  |  |  |

$\dagger$ From Adopted Gammas.

* For absolute intensity per 100 decays, multiply by 0.058 .
\# Total theoretical internal conversion coefficients, calculated using the BrIcc code (2008Ki07) with Frozen orbital approximation based on $\gamma$-ray energies, assigned multipolarities, and mixing ratios, unless otherwise specified.



[^0]:    Continued on next page (footnotes at end of table)

