

**$^{40}\text{K}$   $\varepsilon$  decay ( $1.248 \times 10^9$  y) [1999BeZQ](#), [1999BeZS](#)**

Type	Author	History Citation	Literature Cutoff Date
Full Evaluation	Jun Chen	NDS 140, 1 (2017)	30-Sep-2015

Parent:  $^{40}\text{K}$ :  $E=0$ ;  $J^\pi=4^-$ ;  $T_{1/2}=1.248 \times 10^9$  y 3;  $Q(\varepsilon)=1504.40$  6;  $\% \varepsilon + \% \beta^+$  decay = 10.72 11

$^{40}\text{K}$ - $J^\pi$ : From unique 3rd forbidden  $\beta^-$  spectral shape for decay to  $0^+$  level and L transfer in charge-particle reactions.

$^{40}\text{K}$ - $T_{1/2}$ : From [2004Ko09](#) and [2002Gr01](#); the same value from measurements of specific activity of natural potassium salts using liquid-scintillation counting (LSC) technique. ([2002Gr01](#) reported a value of  $1.248 \times 10^9$  y 2, later adjusted to  $1.248 \times 10^9$  y 3 by [2004Ko09](#) to correct the quoted uncertainty on measured isotopic abundance of  $^{40}\text{K}$ ). Both papers used natural abundance of  $^{40}\text{K}$  as 0.01167% 2 ([1975Ga24](#)). The natural abundance of  $^{40}\text{K}=0.0117\%$  1 (as recommended in the International Union of Pure and Applied Chemistry 70, 217 (1998), based on the measured value of [1975Ga24](#)) would give about four times larger uncertainty on  $T_{1/2}$ . The earlier values of  $1.265 \times 10^9$  y 13 ([1999BeZS](#), [1999BeZQ](#)) based on recomputation of  $1.277 \times 10^9$  y 8 (evaluation by [1973EnVA](#)); and  $1.26 \times 10^9$  y 1 (evaluation by [1990Ho28](#) from 14 different measurements out of a total of 34 measurements listed) are in good agreement. Variation of  $T_{1/2}$  due to environmental conditions has been studied by [2001No10](#), where No significant effect has been reported. Earlier (pre-1977) measurements of partial ( $\beta^-$  and ce) and/or total  $T_{1/2}$  of  $^{40}\text{K}$ : [1977Ce04](#), [1972Go21](#), [1966Fe09](#), [1965Le15](#), [1965Br25](#), [1962Fl05](#), [1961GI07](#), [1960Sa31](#), [1960Eg01](#), [1959Ke26](#), [1957We43](#), [1956Mc20](#), [1955Ba25](#), [1955Ko21](#), [1955Su38](#), [1953Bu58](#), [1950Sa52](#), [1947GI07](#). Another 16 references (from 1931 to 1971) are listed by [1990Ho28](#) and in the 1978 Table of Isotopes ([1978LeZA](#)); but are not present in the NSR database.

$^{40}\text{K}$ - $T_{1/2}$ : @B@0@0@@@ @B@0@1@@@@@1  $T_{1/2}=3.992 \times 10^{16}$  s 40 or  $1.265 \times 10^9$  y 13.

$^{40}\text{K}$ - $Q(\varepsilon)$ : From [2012Wa38](#).

$^{40}\text{K}$ - $\% \varepsilon + \% \beta^+$  decay: deduced by the present evaluator based on  $I\gamma(1460\gamma)/I\beta^- = 0.1195$  14, which is equal to  $I(\varepsilon$  to 1461 level)/ $I\beta^-$ , and  $I(\beta^+)/I(\beta^-) = 1.12 \times 10^{-5}$  14 from evaluation of [1973EnVA](#), and  $\varepsilon/\beta^+$  ( $^{40}\text{K}$  to  $^{40}\text{Ar}$  g.s.) = 45.2 14 (3U theory), with all  $\beta^+$  decay proceeding to  $^{40}\text{Ar}$  ground state. Previously evaluated value by [1999BeZQ](#), [1999BeZS](#) is 0.1086 13 based on the estimation of  $\varepsilon/\beta^+ = 200$  100 for the unique 3rd forbidden branch to the  $^{40}\text{Ar}$  ground state.

**Additional information 2.**

[1999BeZQ](#), [1999BeZS](#): evaluations of  $^{40}\text{K}$  decay.

Measurements: [2014Be25](#), [2013Be06](#), [2004Ko09](#), [2002Gr01](#), [2001No10](#), [1977Ce04](#), [1972Go21](#), [1967Mc10](#), [1966Fe09](#), [1965Le15](#), [1965Br25](#), [1962Fl05](#), [1962En01](#), [1961GI07](#), [1960Sa31](#), [1960Eg01](#), [1959Ke26](#), [1957We43](#), [1956Mc20](#), [1955Ba25](#), [1955Ko21](#), [1955Su38](#), [1953Bu58](#), [1952Fe16](#), [1951Go29](#), [1951De34](#), [1950Sa52](#), [1949Ov01](#), [1948Ev09](#), [1947GI07](#). This list is not complete, see [1978LeZA](#) for several other references that are not present in NSR database.

The decay scheme, which includes the  $\beta^-$  decay to the ground state of  $^{40}\text{Ca}$  and two levels in  $^{40}\text{Ar}$ , is complete since these are the only levels in the daughter nuclides below the respective decay energies.

In principle, the 1460-keV  $\gamma$  ray could be used for energy calibration. However, in a Ge semiconductor detector the apparent  $\gamma$ -ray energy depends on the source-detector configuration and  $^{40}\text{K}$  sources usually consist of a large volume of material, so this  $E\gamma$  is usually not useful. This also means that in most cases the uncertainty in the observed energy is much larger than that given here.

 $^{40}\text{Ar}$  Levels

E(level)	$J^\pi$	$T_{1/2}$	Comments
0	$0^+$	stable	
1460.851 6	$2^+$		$J^\pi$ : from Adopted Levels.

 $\varepsilon, \beta^+$  radiations

E(decay)	E(level)	$I\beta^+$ †	$I\varepsilon$ †	Log $ft$	$I(\varepsilon + \beta^+)$ †	Comments
(43.55 6)	1460.851		10.67 11	11.53 <sup>1u</sup> 1	10.67 11	$\varepsilon\text{K}=0.7609$ 4; $\varepsilon\text{L}=0.2114$ 3; $\varepsilon\text{M}+=0.02771$ 4
(1504.40 6)	0	0.00100 13	0.045 6	21.4 <sup>3u</sup>	0.046 6	av $E\beta=197.325$ 25; $\varepsilon\text{K}=0.5059$ 1; $\varepsilon\text{L}=0.04906$ 1; $\varepsilon\text{M}+=0.007191$ 2
						$I\varepsilon$ : from $I\beta^+$ (to $^{40}\text{Ar}$ g.s.)/ $I\beta^- = 1.12 \times 10^{-5}$ 14 in evaluation of <a href="#">1973EnVA</a> and adopted $\%I\beta^- = 89.28$ 11, with $\varepsilon/\beta^+$ ( $^{40}\text{K}$ to $^{40}\text{Ar}$ g.s.) = 45.2 14 (3U)

Continued on next page (footnotes at end of table)

**$^{40}\text{K}$   $\epsilon$  decay ( $1.248 \times 10^9$  y) [1999BeZQ,1999BeZS](#) (continued)**

$\epsilon, \beta^+$  radiations (continued)

<u>E(decay)</u>	<u>E(level)</u>	<u>Comments</u>
		theory). Log <i>ft</i> : from private communication from R. B. Firestone; see also <a href="#">1970Wa11</a> . <a href="#">Additional information 3</a> .

† Absolute intensity per 100 decays.

$\gamma(^{40}\text{Ar})$

$I_\gamma$  normalization:  $I_\gamma(1460 \gamma)$  is from the measured  $\gamma/\beta^-$  ratio (evaluated in [1973EnVA](#)), which can be obtained from  $I(\epsilon, 1460)/(1+\alpha+\text{IPFC})$ .  $\alpha(1460)=2.5 \times 10^{-5}$  and  $\text{IPFC}=7.3 \times 10^{-5}$ , so the correction for these is 0.01% and is completely negligible compared to the 1% uncertainty in  $I(\epsilon, 1460)$ .

<u><math>E_\gamma</math></u>	<u><math>I_\gamma^\dagger</math></u>	<u><math>E_i(\text{level})</math></u>	<u><math>J_i^\pi</math></u>	<u><math>E_f</math></u>	<u><math>J_f^\pi</math></u>	<u>Mult.</u>	<u><math>\alpha^\ddagger</math></u>	<u>Comments</u>
1460.820 5	10.66 13	1460.851	2 <sup>+</sup>	0	0 <sup>+</sup>	E2	$2.95 \times 10^{-5}$ 9	$E_\gamma$ : the evaluator has re-scaled the original values in <a href="#">1979He13</a> using the new calibration standards in <a href="#">2000He14</a> . Others: 1460.75 6 ( <a href="#">1967Ki10</a> ), 1460.95 7 ( <a href="#">1970Ja15</a> ). $I_\gamma$ : $I_\gamma(1460)=I(\epsilon, 1460)/(1+\alpha+\text{IPFC})=10.67$ 11/1.000102 5. <a href="#">Additional information 4</a> .

† Absolute intensity per 100 decays.

‡ Total theoretical internal conversion coefficients, calculated using the BrIcc code ([2008Ki07](#)) with Frozen orbital approximation based on  $\gamma$ -ray energies, assigned multipolarities, and mixing ratios, unless otherwise specified.

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Decay Scheme

Intensities:  $I_{(\gamma+ce)}$  per 100 parent decays

