# Multilevel Breit-Wigner (MLBW) elastic angular distributions 

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## Leakage from a critical assembly depends greatly on angular dists.

- A forward peaked distribution lets more particles escape from collision events on boundaries than an isotropic one
- Most evident in small critical systems



## cea <br> Shielding benchmark

## ASPIS benchmark (iron)

Total Monte Carlo results obtained with the TRIPOLI code shows that above 20 cm , the flux attenuation is dominated by higher order polynomial coefficient uncertainties



from $G$. Noguere

## What angular distributions are available in the resonance region?

| Format | Number <br> occurrences | Number with angular <br> distributions enabled |
| :---: | :---: | :---: |
| scattering radius only | 66 | 0 |
| SLBW | 8 | 0 |
| MLBW | 270 | 0 |
| Reich Moore (LRF=3) | 54 | 46 |
| R Matrix Limited (LRF=7) | 1 | 0 |
| URR only | 26 | 0 |

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| Reich | 0 |  |
| Moore (LRF=: can we salvage these guys? |  |  |
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## MLBW Approximation to the R matrix

- R matrix theory gets us the collision matrix (c=channel index)

$$
U_{c c^{\prime}}=e^{-i\left(\varphi_{c}+\varphi_{c^{\prime}}\right)}\left(\delta_{c c^{\prime}}+i \sum_{\lambda, \mu} \Gamma_{\lambda c}^{1 / 2} A_{\lambda \mu} \Gamma_{\mu c^{\prime}}^{1 / 2}\right)
$$

- in terms of the level matrix $A$ :

$$
\left(\mathbf{A}^{-1}\right)_{\lambda \mu}=\left(E_{\lambda}-E\right) \delta_{\lambda \mu}-\sum_{c} \gamma_{\lambda c} L_{c}^{0} \gamma_{\mu c}
$$

- Here, $E_{\lambda}$ is the resonance energy, $y_{\lambda c}$ is the reduced width and $L^{0}{ }_{c}$ is essentially the hard sphere shift and penetrability


## MLBW Approximation to the R matrix

- With just the scattering matrix, we can compute the total cross section

$$
\sigma_{c} \equiv \sum_{c^{\prime}} \sigma_{c c^{\prime}}=2 \pi \lambda_{c}^{2}\left(1-\Re U_{c c}\right)
$$

- and the cross section to channel $c^{\prime}$

$$
\sigma_{c c^{\prime}}=\pi \lambda_{c}^{2} g_{c}\left|\delta_{c c^{\prime}}-U_{c c^{\prime}}\right|^{2}
$$

## MLBW Approximation to the R matrix

- MLBW throws away off diagonal parts of the level matrix and renames things into the widths we are familiar with

$$
\begin{aligned}
\left(\mathbf{A}^{-1}\right)_{\lambda \mu} & =\left(E_{\lambda}-E-\sum_{c} L_{c}^{0} \gamma_{\mu c}^{2}\right) \delta_{\lambda \mu} \\
& \equiv\left(E_{\lambda}+\Delta_{\lambda}-E-i \Gamma_{\lambda} / 2\right) \delta_{\lambda \mu}
\end{aligned}
$$

- Result is that collision matrix no longer unitary; we can measure this with Frobenius norm:

$$
\|U\|_{F}=\sqrt{\sum_{c c^{\prime}} U_{c c^{\prime}}^{*} U_{c^{\prime} c}}
$$

If $U$ is unitary, $\|U\|_{F}=\sqrt{\operatorname{dim}(U)}$

## So MLBW is not unitary. Is this a problem?



## ENDF’s MLBW != R matrix MLBW

- ENDF MLBW uses SLBW formulas for non-elastic cross sections


## Bad

- ENDF overrides use of potential scattering radius in phase-factor of scattering matrix


## Meh

## ENDF's MLBW capture is clearly different



## This naturally impacts the total too



## ENDF MLBW elastic is the same as the $\mathbf{R}$ matrix MLBW elastic

Elastic


## Blatt-Beidenharn shows us how to make angular distributions

According to Blatt-Biedenharn formalism, we only need a collision matrix in order to compute an angular distribution:

$$
\frac{d \sigma_{\alpha, \alpha^{\prime}}(E)}{d \Omega}=\frac{1}{k^{2}(2 i+1)(2 I+1)} \sum_{s, s^{\prime}} \sum_{L=0}^{\infty} B_{L}\left(\alpha s, \alpha^{\prime} s^{\prime} ; E\right) P_{L}(\mu)
$$

where

$$
\begin{aligned}
& B_{L}\left(\alpha s, \alpha^{\prime} s^{\prime} ; E\right)= \\
& \quad \frac{(-)^{s-s^{\prime}}}{4} \sum_{c_{1}=\left\{\alpha \ell_{1} s_{1} J_{1}\right\}} \sum_{c_{1}^{\prime}=\left\{\alpha^{\prime} \ell_{1}^{\prime} s_{1}^{\prime} J_{1}^{\prime}\right\}} \sum_{c_{2}=\left\{\alpha \ell_{2} s_{2} J_{2}\right\}} \sum_{c_{2}^{\prime}=\left\{\alpha^{\prime} \ell_{2}^{\prime} s_{2}^{\prime} J_{2}^{\prime}\right\}} \bar{Z}\left(\ell_{1} J_{1} \ell_{2} J_{2} s L\right) \bar{Z}\left(\ell_{1}^{\prime} J_{1} \ell_{2}^{\prime} J_{2} s^{\prime} L\right) \\
& \quad \times \delta_{s s_{1}} \delta_{s^{\prime} s_{1}^{\prime}} \delta_{J_{1} J_{1}^{\prime} \delta_{s s_{2}} \delta_{s^{\prime} s_{2}^{\prime}}^{\prime} J_{J_{2} J_{2}^{\prime}}\left(\delta_{c_{1} c_{1}^{\prime}}-U_{c_{1} c_{1}^{\prime}}(E)\right)^{*}\left(\delta_{c_{2} c_{2}^{\prime}}-U_{c_{2} c_{2}^{\prime}}(E)\right)}
\end{aligned}
$$

## We can get MLBW angular distributions!!!

- For the MLBW format, we have such a thing:

$$
U_{c c^{\prime}}=e^{-i\left(\varphi_{c}+\varphi_{c^{\prime}}\right)}\left(\delta_{c c^{\prime}}+\sum_{\lambda} \frac{i \Gamma_{\lambda c}^{1 / 2} \Gamma_{\lambda c^{\prime}}^{1 / 2}}{E_{\lambda}+\Delta_{\lambda}-E-i \Gamma_{\lambda} / 2}\right)
$$

- The MLBW formalism in ENDF is "broken", but it's elastic channel and the collision matrix are not

OK then, what does it look like?

## The angular distribution directly from ENDF MLBW parameters



## Try it on a more popular isotope: ${ }^{90} \mathrm{Zr}$ distribution of level widths



## Try it on a more popular isotope: ${ }^{90} \mathbf{Z r}$ angular distribution



## Try it on a more popular isotope: ${ }^{90} \mathbf{Z r}$ zoom in on angular distribution



## What next?

- Finish coding hooks in Fudge
- Get the distributions into an ENDF file (e.g. 90Zr)
- Compare to ${ }^{\text {nat } Z r} \bar{\mu}$
- Try it out in a benchmark

