History of Pu-239 fission cross section measurements

Fredrik Tovesson Los Alamos National Laboratory



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Outline

Experimental approaches

- Detectors
- Neutron sources
- Flux monitoring
- Data analysis
- Uncertainties

Experimental data in EXFOR

- Pre-1970
- Absolute measurements
- Relative measurements up to 20 MeV
- Spallation source measurements



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Detection methods

- Ionization chambers
 - Fritsch-gridded
 - Parallel plate
 - PPAC
- Surface barrier
- Scintillators
- Track-etch detectors









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Housing – Aluminum OD 1.740" .040" Thick

> Flex Circuit with Spring Contacts*

FEFEFEF

Index Feature in Housing

Target Cassette

Pull Tab Rear Window

Connector to

Amp Circuits

Ring -Aluminum

Neutron production

Quasi mono-energetic neutrons

- Accelerators: Van de Graaff (Geel, Ohio State U.), cyclotron (UU)
- Li(p,n): 0 3.8 MeV (above 3.8 the Li(p,n') reaction contributes)
- T(p,n): 1.0 5.0
- D(d,n): 3.5 10.0 MeV (deuterium break-up above 7 MeV)
- T(d,n): 14.1-22.0 MeV
- White neutron sources
 - Photonuclear neutron production (ORELA, GELINA).
 - Electron beam (100 MeV) on uranium target (Geel). Energy range from sub-thermal to 20 MeV.
 - Spallation (LANSCE, PNPI, nTOF)
 - Proton beam on heavy target, such as lead or tungsten
 - Neutron range sub-thermal to hundreds of MeV's



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Flux measurements

- **Relative cross section measurements**
 - U-235(n,f). ٠
 - Standard evaluation
 - ~1% uncertainty in fast region
 - Many systematic uncertainties in the fission counting are reduced
 - H(n,n)H ٠
 - The proton recoil in neutron scattering on hydrogen is detected
 - Uncertainty is fraction of percent
 - Detected with different system than the fission event = larger systematic uncertainties
 - $B(n,\alpha)$, $Li(n, \alpha)$ ٠
 - Standards in the low-energy range (below 100 keV)
- Absolute cross section measurements
 - Associated particle method
 - In the neutron producing reaction the associated particle is detected,



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Cross section analysis





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Neutron background





Detection efficiency



- Fission fragment absorption in sample
 - Thin (<200 ug/cm2) samples are typically used. Still, some 2-3% of events are absorbed
- Energy straggling



Straggling causes a tail of the energy distribution, and some events fall below detection threshold

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Dead time corrections



P. B. Coates, Rev. Sci. Instrum. 63, 1992



 p_i = probability (per T0) to have event in bin *i*

 N_i = measured events in bin *i*

 N_p = number of T0's

D = deadtime in bins

- The dead-time correction is exact at constant event rates. We correct every 10 seconds worth of data, assuming stable beam on this time scale.
- Hardware scalers measures the integral live-time. The only input parameter D is fine tuned until there is perfect agreement with scalers.

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Fission identification



Background events

- Radioactive decay
- Neutron-induced chargedparticles
- Gammas
- Rejection criteria
 - Total energy deposited
 - TOF



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Mass determination



Spectrum from alpha counting of Pu-244 samples used for fission cross section measurement by Staples and Morley

- Samples are typically counted relative to standards. Pu-239 samples are alpha- or gamma counted.
- Enriched samples are needed for fission cross sections. Pu-239 is available with higher than 99% enrichment. Contamination levels can be determined using alpha spectroscopy.



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Sample and beam profile non-uniformities



- Sample uniformity
 - Highly dependent on method of preparation
 - Evaporated samples can be uniform to within 1-2%
 - Can be measured using counting with mask
- Beam profile
 - Imaging techniques used to determine profile
- Typically not energy differential
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Covariance of experimental uncertainties

Elements in the covariance matrix are given by

$$C_{ij} = R_i R_j \left[\sigma_s^2 \delta_{ij} + d_i d_j + b_i b_j + u_i u_j + p_i p_j + \sigma_n^2 \right]$$

Where

 $\begin{array}{l} \mathsf{R}_{i,j} = \text{measured ratio value} \\ \sigma_s = \text{relative statistical uncertainty} \\ \mathsf{d}_{i,j} = \text{relative uncertainty due to dead-time corrections} \\ \mathsf{b}_{i,j} = \text{relative uncertainty due to background corrections} \\ \mathsf{u}_{i,j} = \text{relative uncertainty due to U-233 contamination corrections} \\ \mathsf{p}_{i,j} = \text{relative uncertainty due to Pu-239 contamination corrections} \\ \sigma_s = \text{normalization uncertainty} \end{array}$



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Experimental work pre-1970



- Generally higher cross sections than current evaluation
- Large discrepancies above 10 MeV

Fairly large experimental uncertainties, or non at all
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Absolute measurements

Author	Year	Energy (MeV), points	Neut. Prod.	Detector	Flux mon.
I. D. Alkhazov et al.	1986	4.8-18.8, 3	D-D, D-T	lon. Chamber	Assoc. part, α
Xianjian et al.	1982	1.0-5.8,16	P-T, D-D	lon. Chamber	H(n,n)H, p
M.C.Davis et al.	1978	0.14-0.96, 4	Na-Be, La-Be, Na-D, and Ga-D	track-etch detectors	Calibrated N- source
K. Kari	1978	0.99-20.9, 168	?, cyclotron, TOF	Gas scint.	H(n,n)H, p
I. Szabo et al.	1976	0.035-5.53, 54	P-Li, P-T,D-D	Frisch	Assoc. part, H(n,n)H



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Absolute measurements





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Relative measurements in the fast region



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Measurements at spallation sources



- Lisowski et al. performed measurements in late 1980's at LANSCE
- Shcherbakov et al. made similar measurements at PNPI in 1990's. Difference: flight path length, pulse spacing.
- New measurements at LANSCE performed in recent years



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Summary

- Large data base of Pu-239 energy differential fission cross sections.
- Most data set were collected with the same type of detectors, neutron sources, samples.
- The uncertainties in all individual measurements are >1%. More realistically >2%.
- Fairly large discrepancies exist above 10 MeV.
- There is plenty of measurements relative to U-235. An absolute, or relative to H(n,n), would be of significantly more value.



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