

CURRENT INCONSISTENCIES IN ^{238}Pu , $^{242,243}\text{Am}$ AND ^{242}Cm EVALUATIONS AND IMPACT ON UNCERTAINTIES

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- Improvements of the nuclear reaction modeling and nuclear parameter systematic for major actinides ^{232}Th , ^{233}U , ^{235}U , ^{238}U , ^{239}Pu
- Sound basis for critical assessment of $\Sigma(n, F)$, $\Sigma(n, \gamma)$, $\Sigma(n, n')$, $\Sigma(n, xn)$ of MA
- Essential «tweaking » of major actinides data files to yield $K\text{-eff} \approx 1$
- A number of biases is imposed into capture, inelastic css, PFNS, etc.
- To avoid substituting model deficiency uncertainties for MA by enlarging the δx_i of x_i nuclear model parameters !!!
- Uncertainty estimation of CSS & PFNS of MA should be preceded by neutron data re-evaluation
- Otherwise, for MA(Np, Pu, Am, Cm) the artificially large $\Delta\Sigma$ - unavoidable

“STRUCTURED” CORRELATIONS

- Close interplay between EXP.(D\$I)& MOD.
 - Correlation of ARP &CSS
 - Correlation of nuclear model parameters
 - Correlation between reaction channels
 - Correlation between (Z,A) & $(Z\pm n, A\pm m)$
 - Correlation between $\sigma(n, F)$ & v_p & PFNS

SCOPE

- 238-Pu
- ENDF/B-VII.0 (ENDF/B-VI)
- JENDL-3.3
- Maslov et al., 1998 (<http://www-nds.iaea.org>)
- 241-Am [Maslov et al. 1998] in JENDL-3.3
- JEFF-3.1 , ENDF/B-VII.0
- 242g-Am [Maslov et al. 1998] in JENDL-3.3, JEFF-3.1
 - ENDF/B-VII.0
- 242m-Am [Maslov et al. 1998] in JENDL-3.3, JEFF-3.1
 - ENDF/B-VII.0
- 243-Am- [Maslov et al. 1998] in JENDL-3.3, JEFF-3.1
 - ENDF/B-VII.0
- 242-Cm
 - ENDF/B-VII.0 (ENDF/B-VII.0)
 - JENDL-3.3
 - BROND
- 243-Cm- Maslov et al. 1998 in JEFF-3.1, JENDL-3.3,ENDF/B-VII
- 245-Cm- Maslov et al. 1998 in JEFF-3.1, JENDL-3.3, ENDF/B-VII
- 246-Cm- Maslov et al. 1998 in JEFF-3.1, JENDL-3.3, ENDF/B-VII

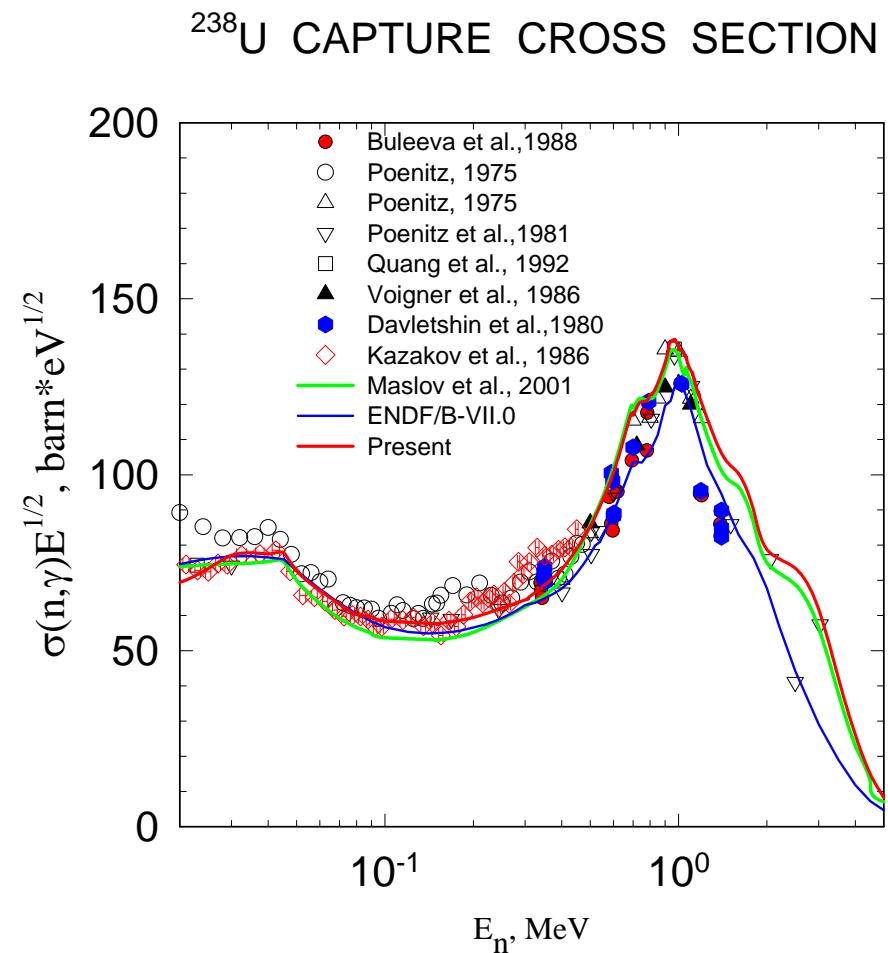
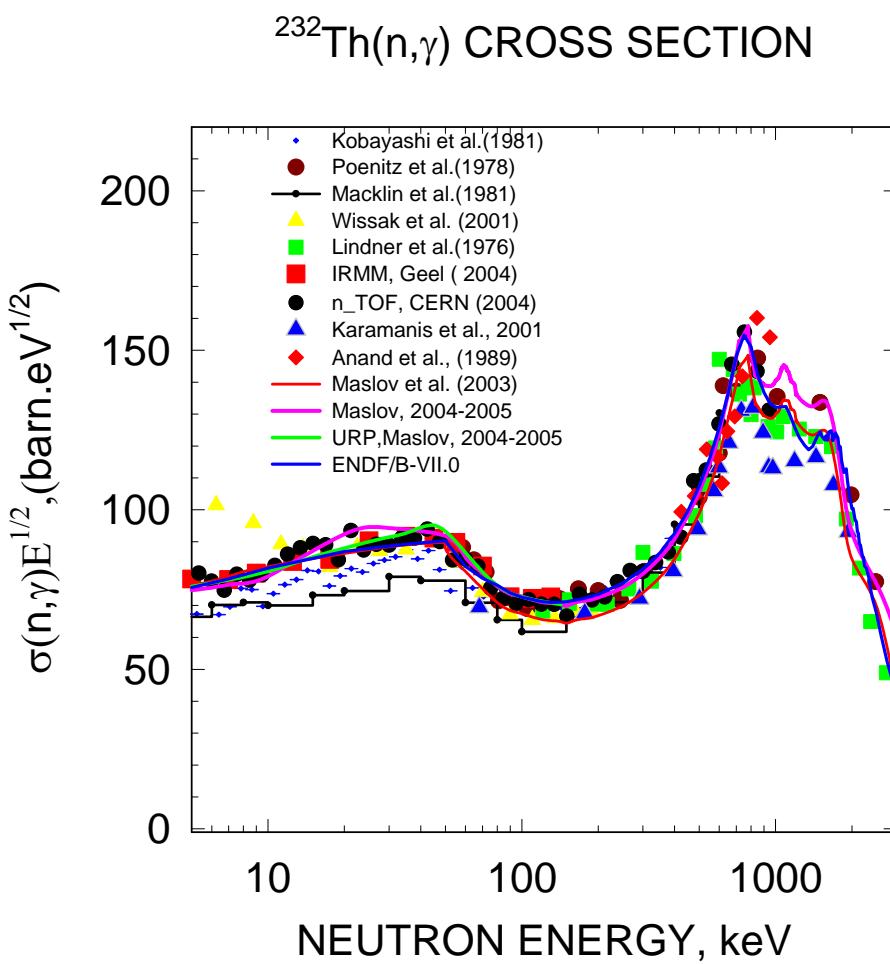
CAPTURE –model deficiency or ARR+OMP uncertainties

^{233}U , ^{235}U , ^{239}Pu

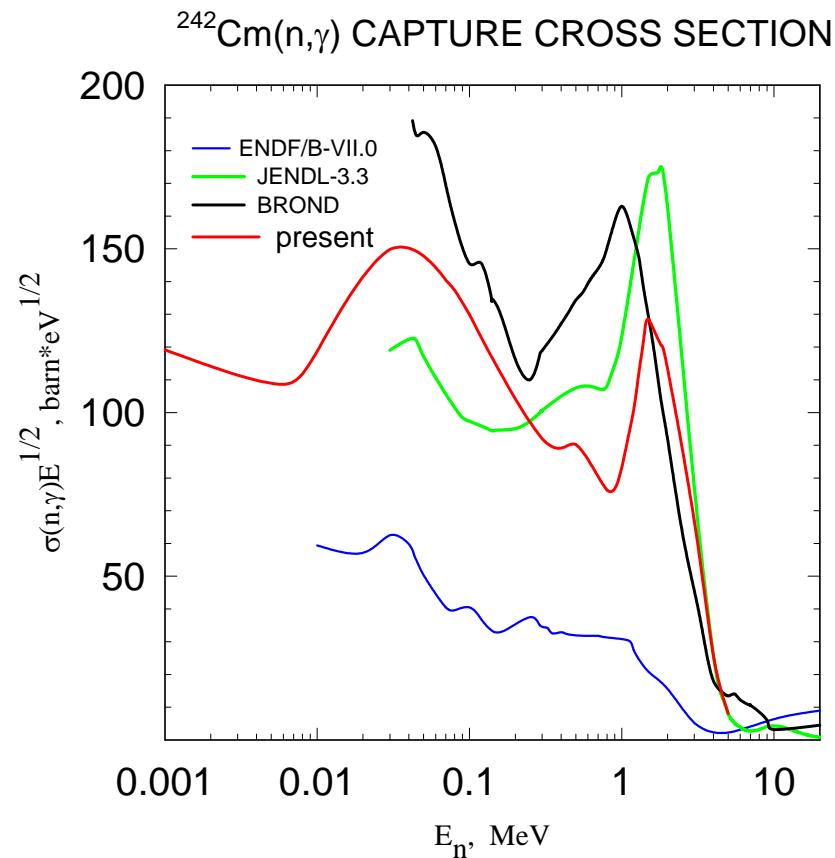
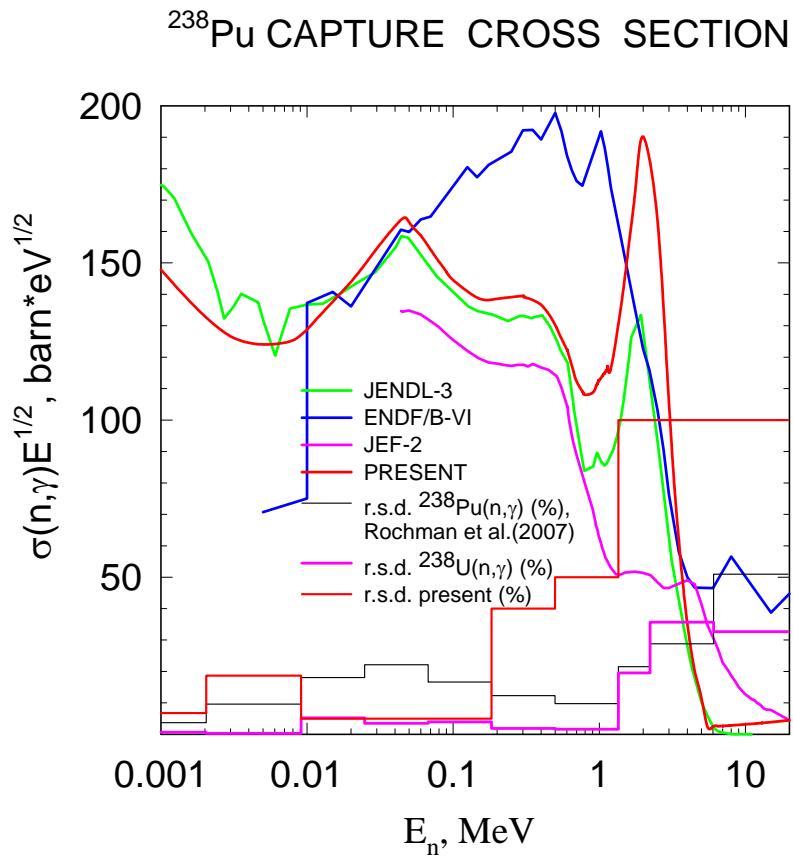
^{232}Th , ^{238}U

^{238}Pu , ^{242}Cm

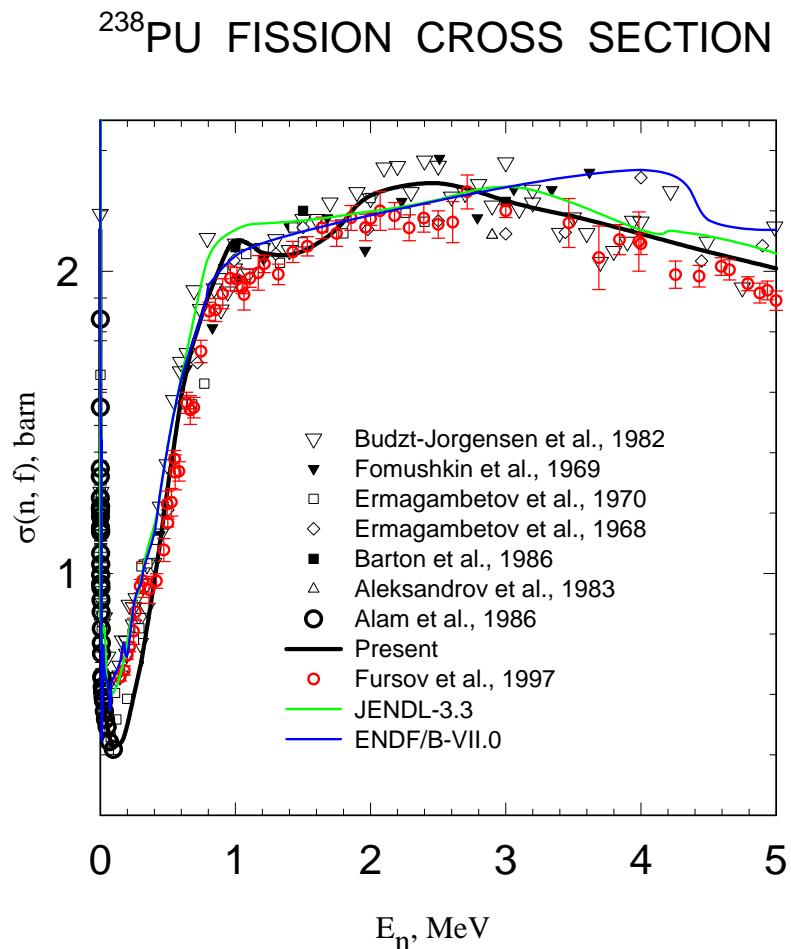
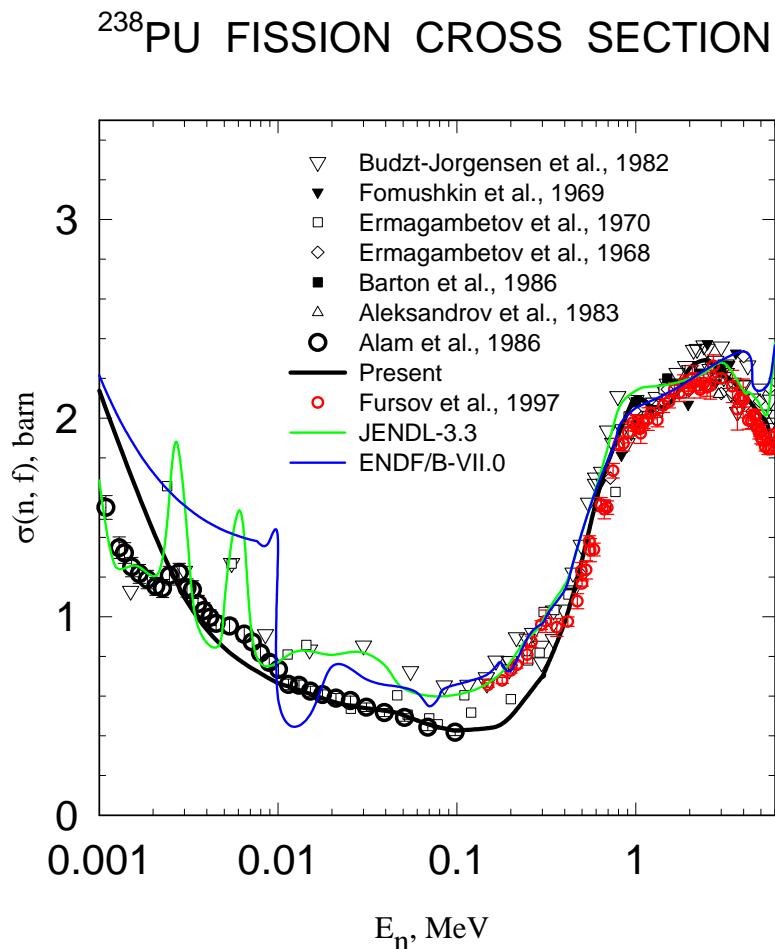
“Tweaking “ consequences- “biases” revealed



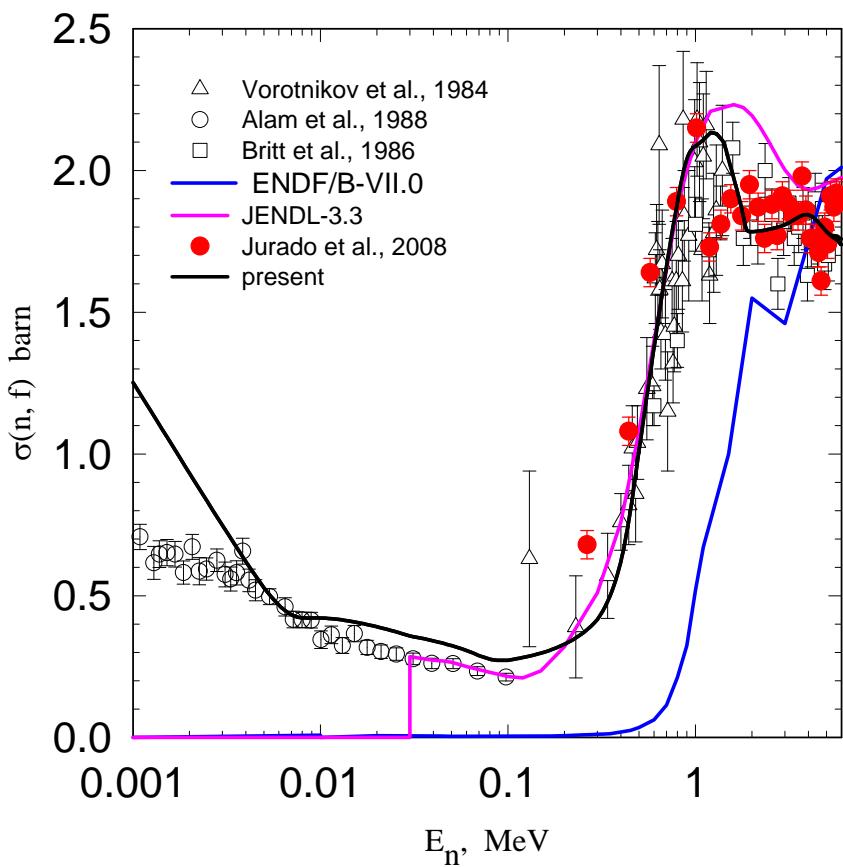
“Tweaking “impossible- ARP+ $\sigma(n,f)$ fit



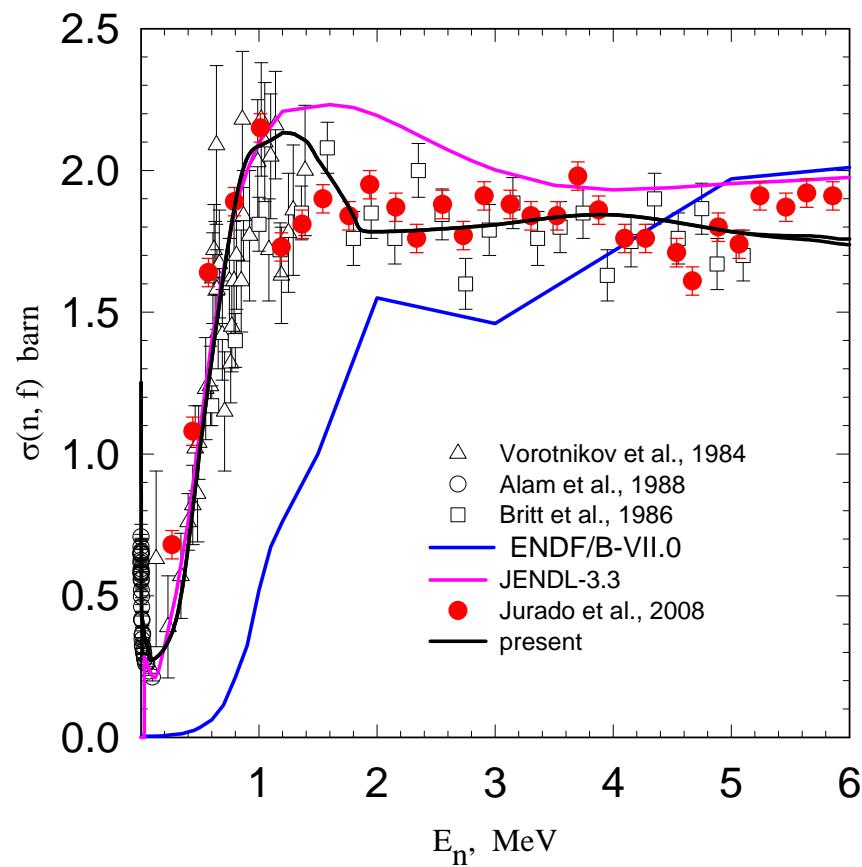
$^{238}\text{Pu}(n, f)$

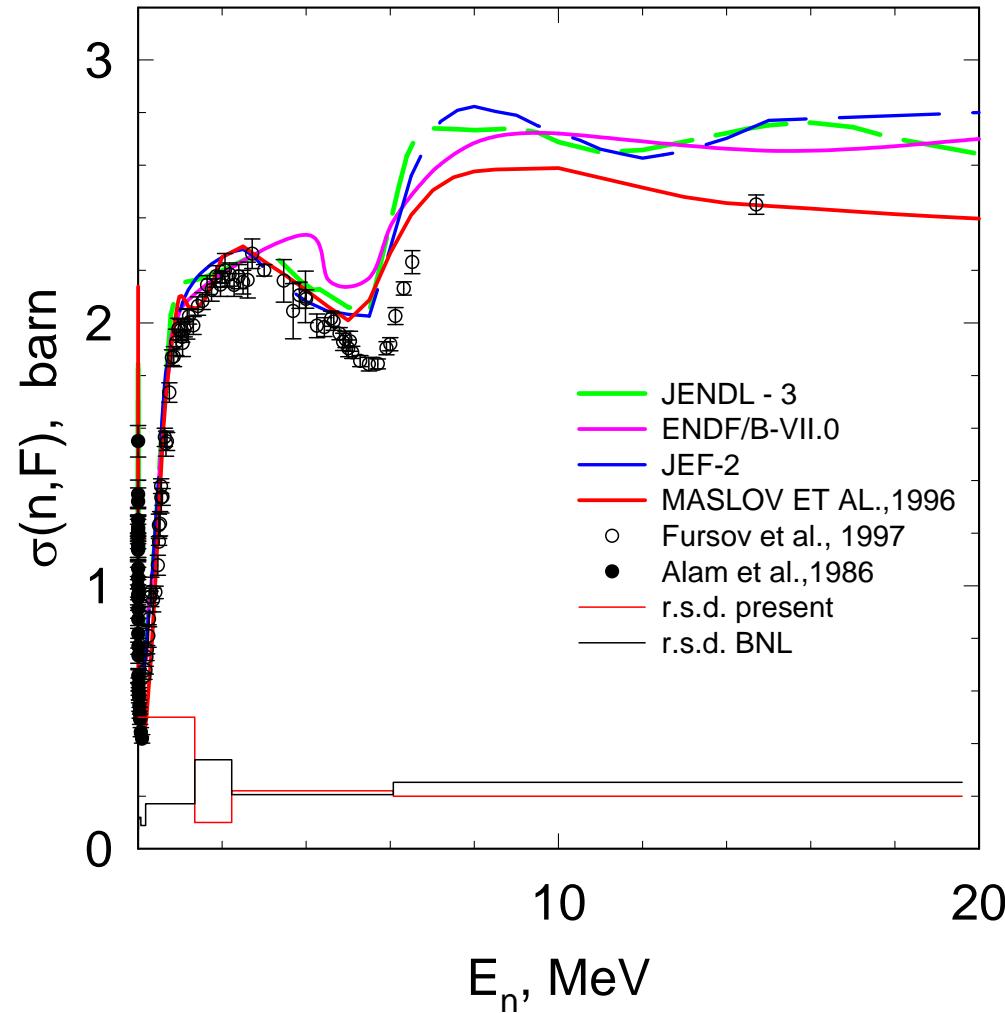


^{242}Cm FISSION CROSS SECTION



^{242}Cm FISSION CROSS SECTION



^{238}Pu FISSION CROSS SECTION

FISSION -model deficiency or Fiss.param. uncertainties

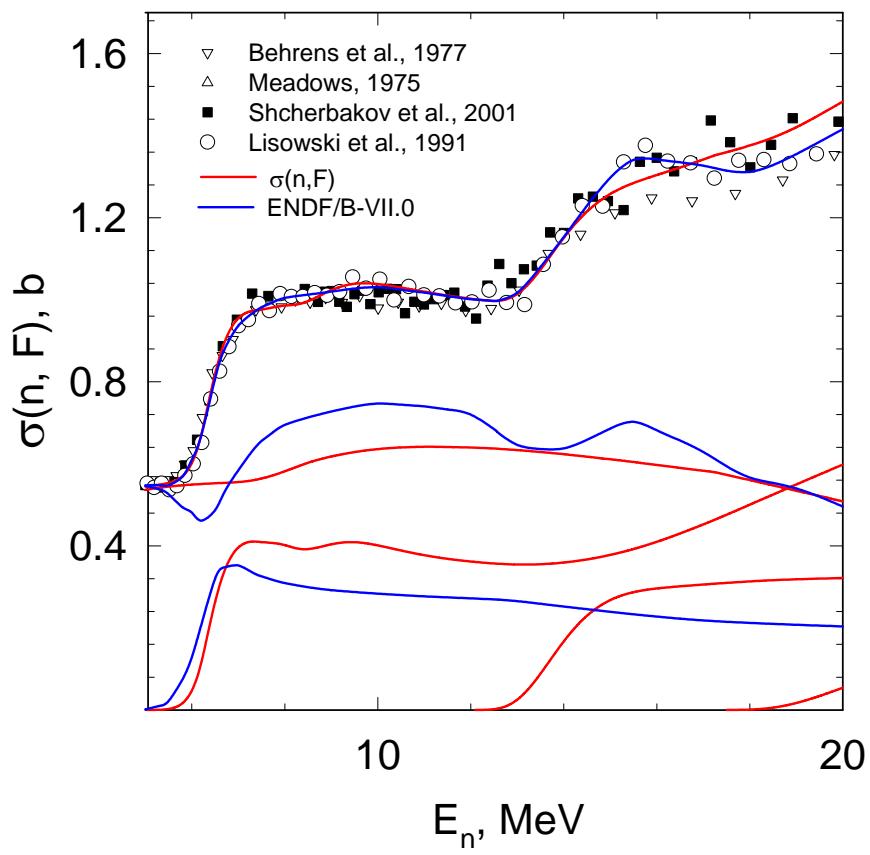
^{235}U , ^{238}U , ^{232}Th

^{237}U

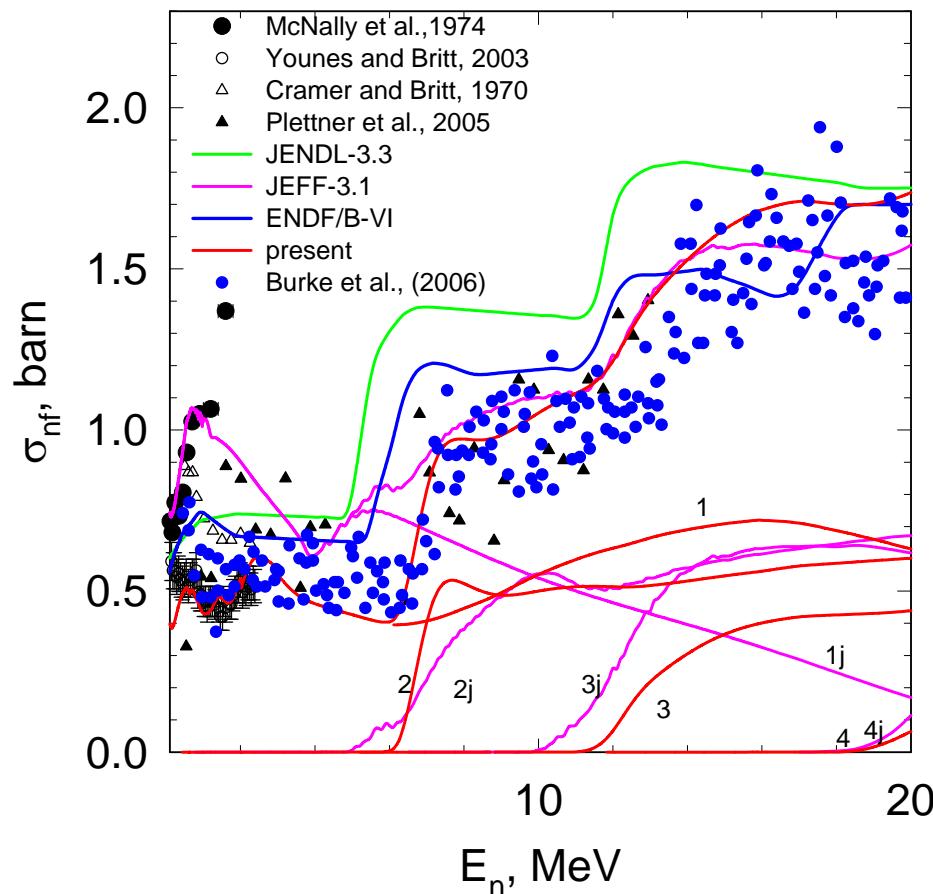
^{238}Pu , ^{242}Cm , ^{241}Am , ^{243}Am

$^{238}\text{U}(n,F)$

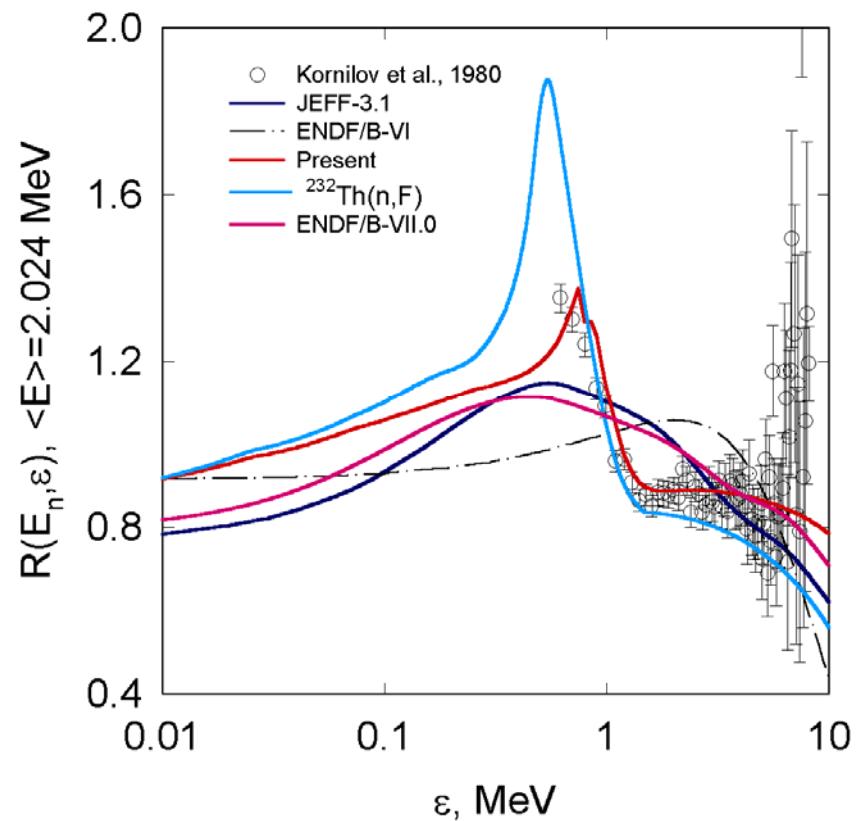
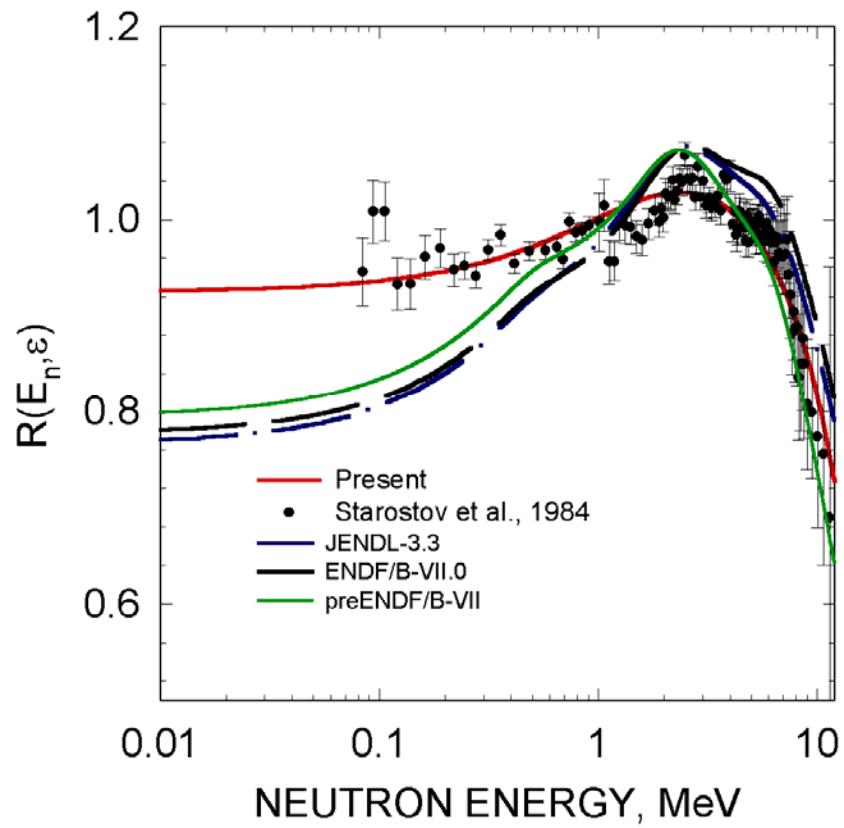
^{238}U FISSION CROSS SECTION

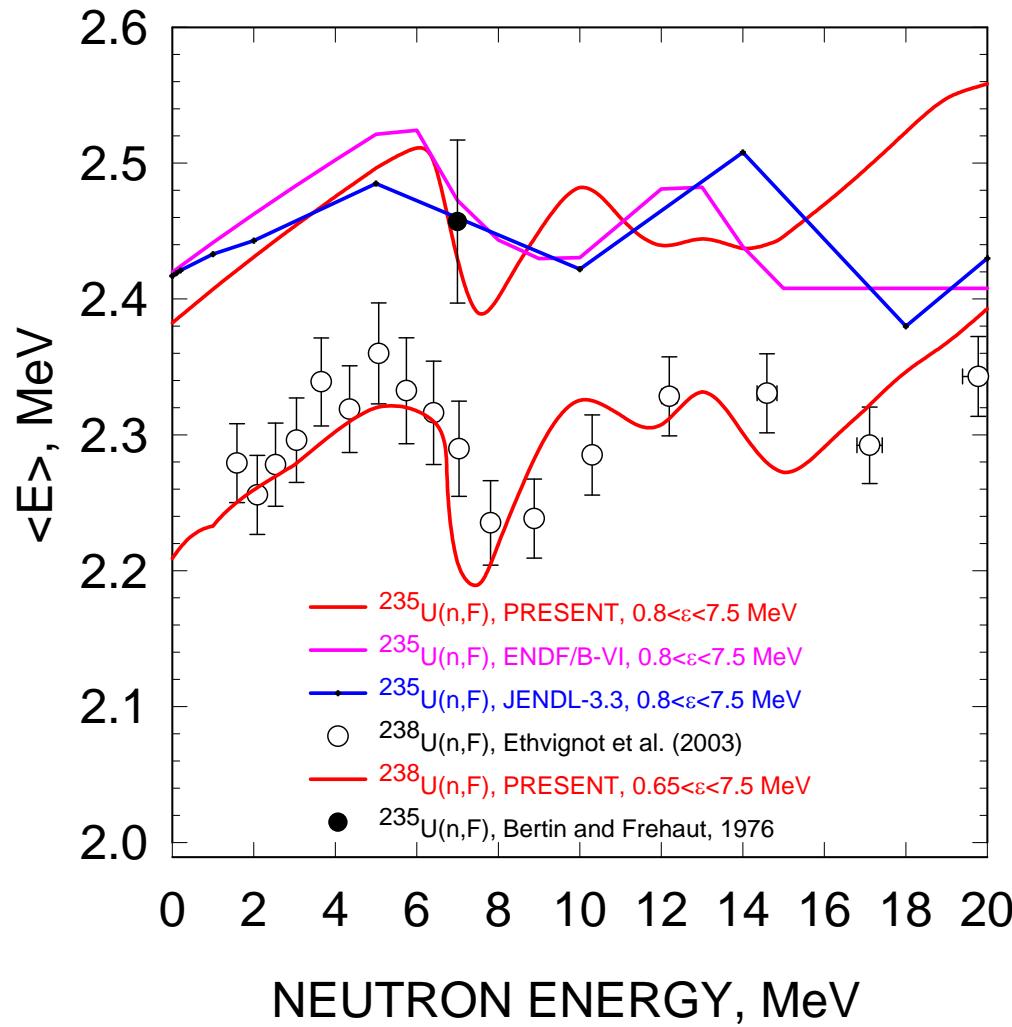


$^{237}\text{U}(n,F)$



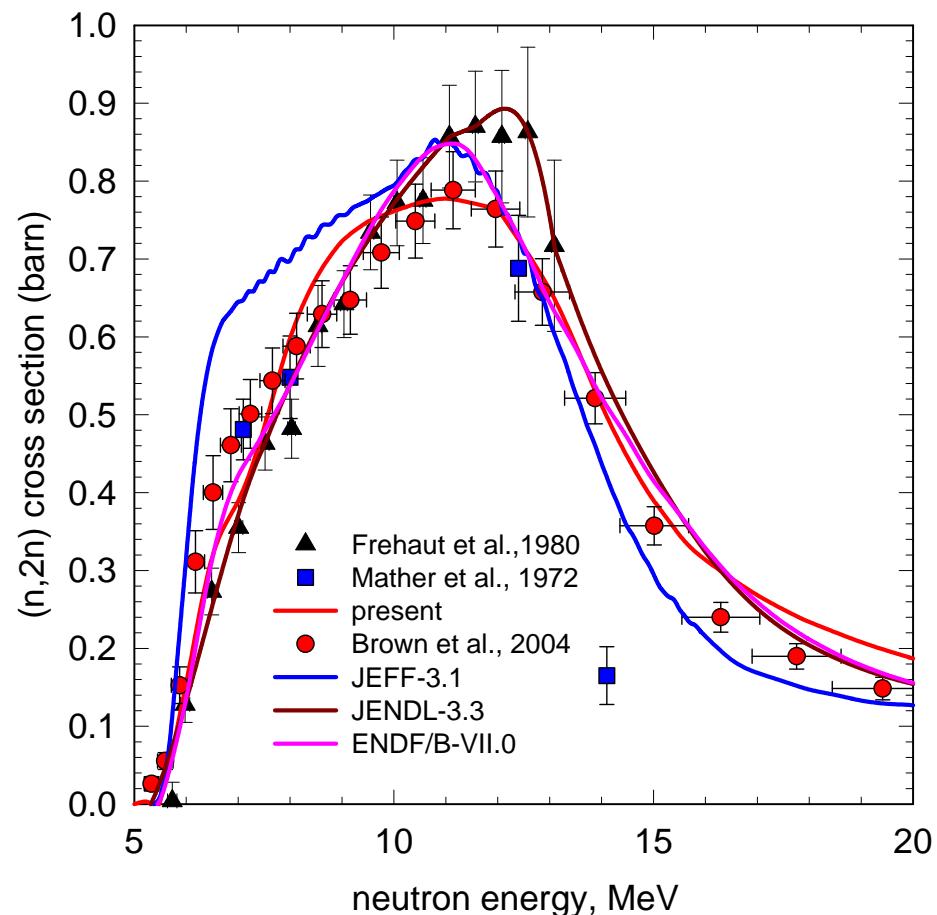
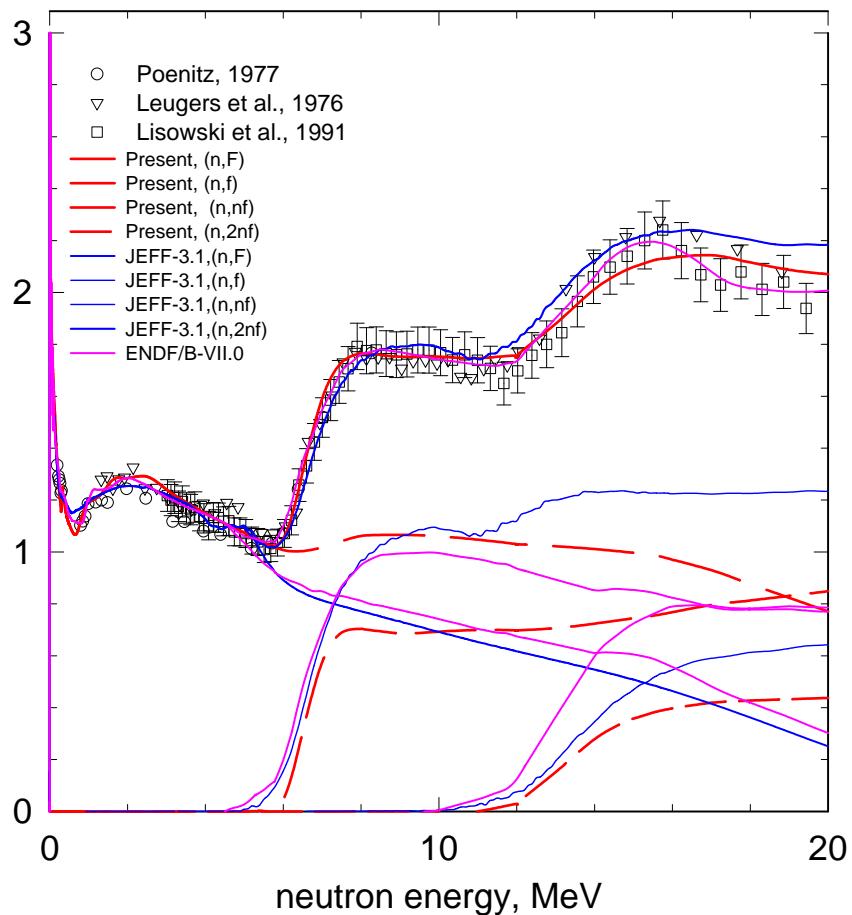
$E_n = \text{thermal}$, $^{238}\text{U}(n, f)$, PFNS, $E_n = 7 \text{ MeV}$



^{235}U & ^{238}U : AVERAGE ENERGY OF PFNS

$^{235}\text{U}(n,\text{F})$

&

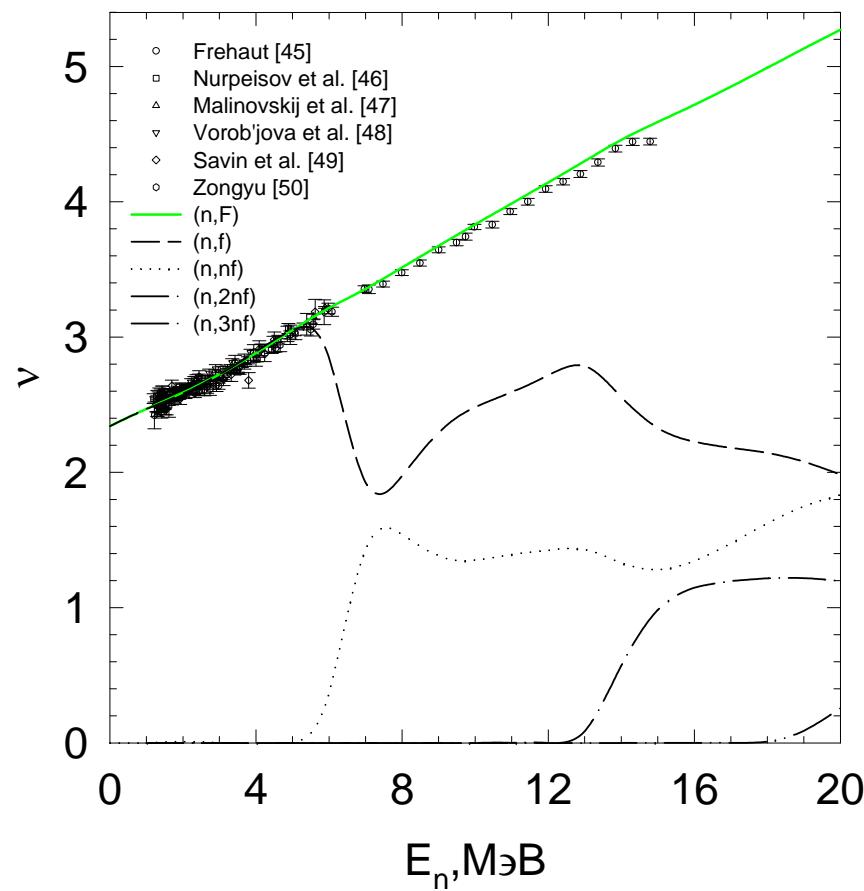
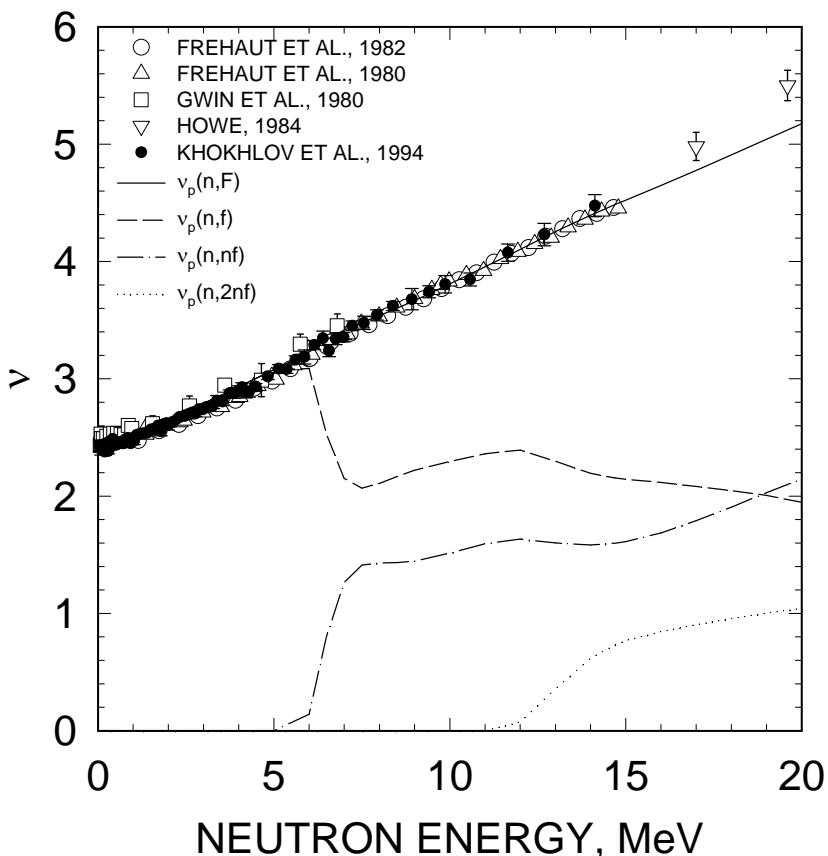
 $^{235}\text{U}(n,2n)$ 

$^{235}\text{U}(\text{n},\text{F})$

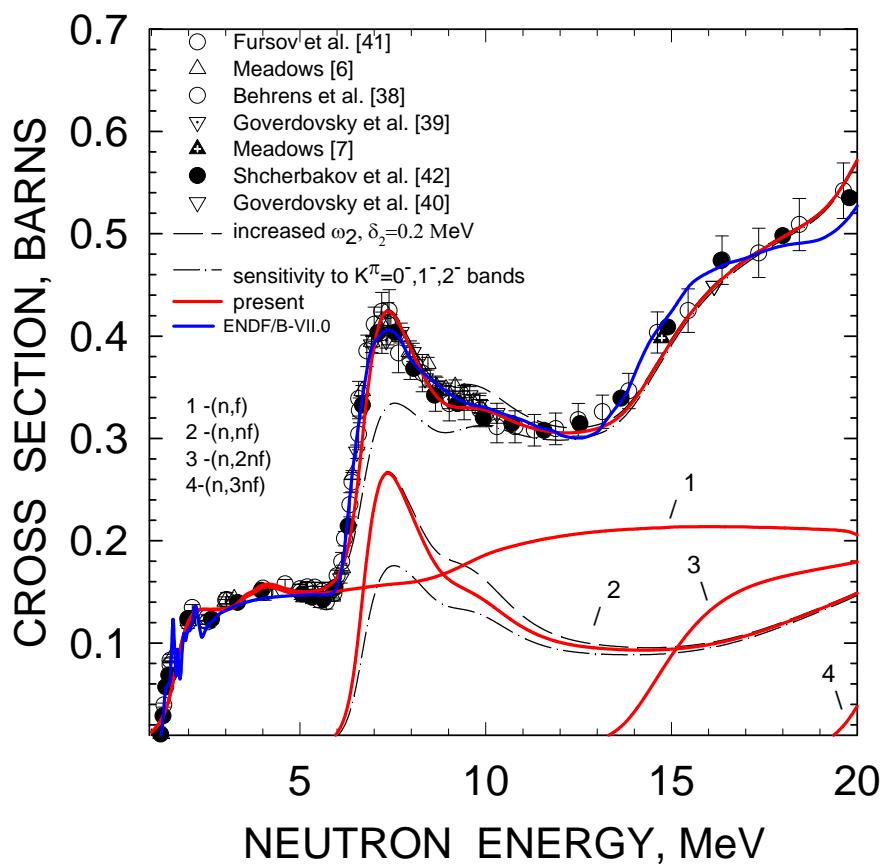
ν of PFN

$^{238}\text{U}(\text{n},\text{F})$

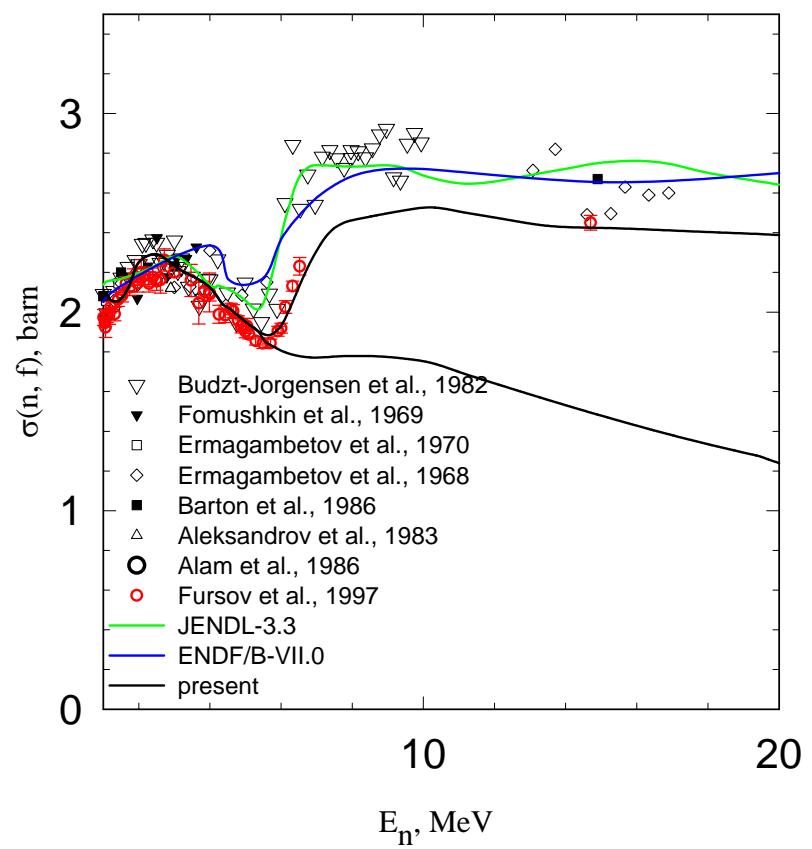
$^{235}\text{U}(\text{n},\text{F})$, NEUTRON MULTIPLICITY



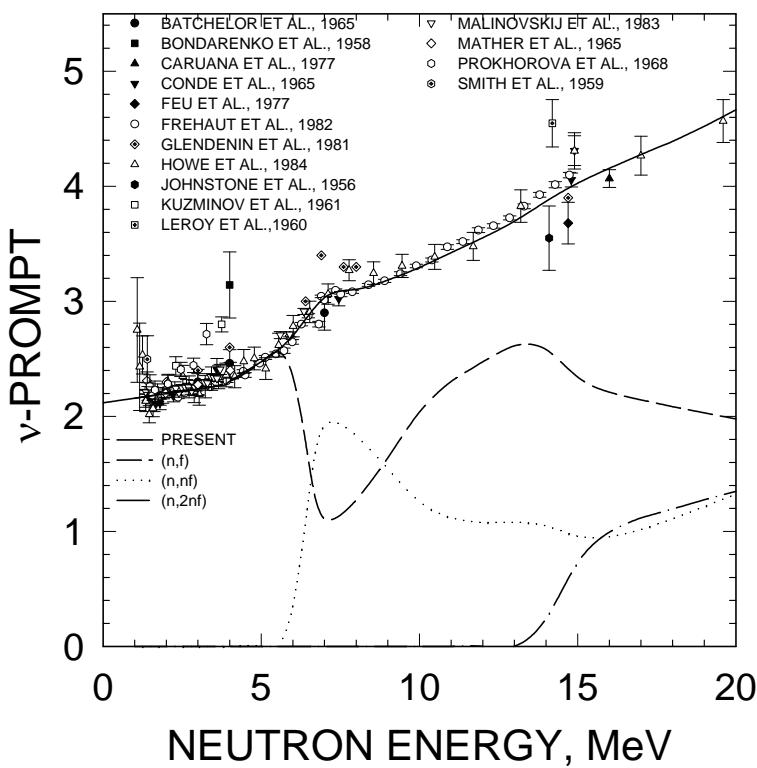
^{232}Th FISSION CROSS SECTION



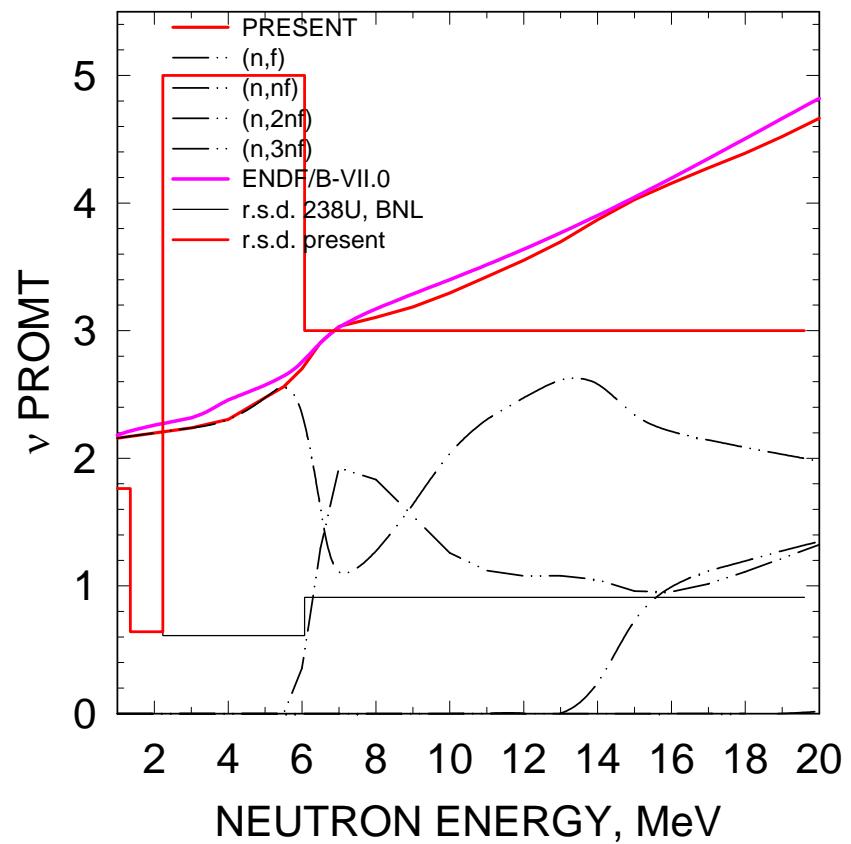
^{238}Pu FISSION CROSS SECTION



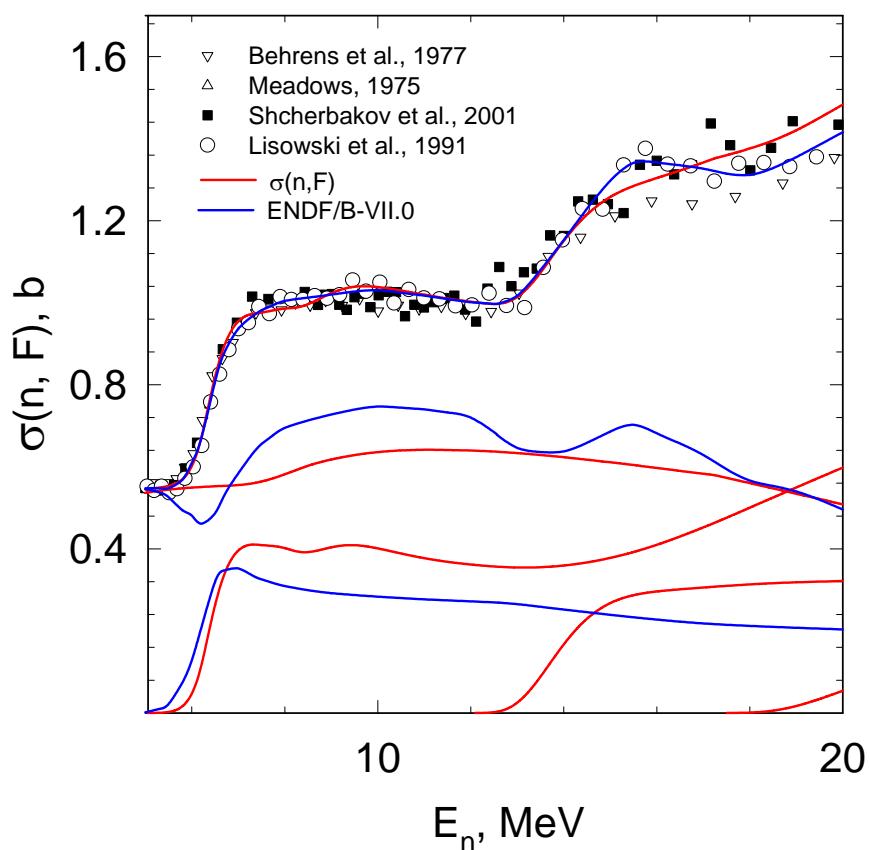
^{232}Th , NEUTRON MULTIPLICITY



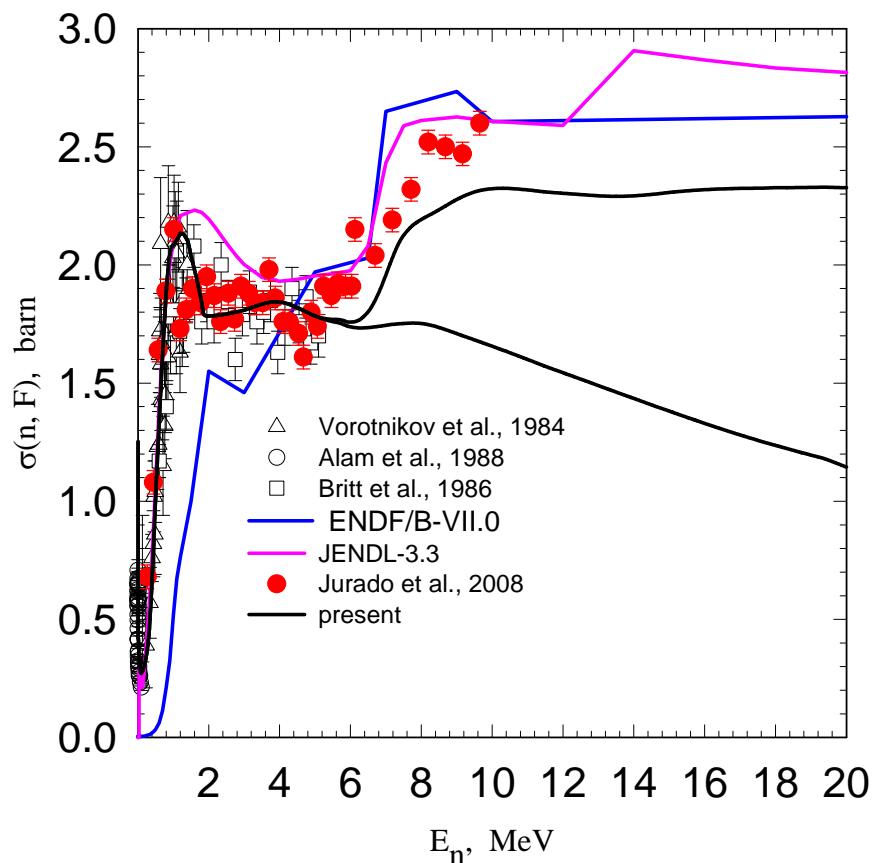
$^{232}\text{Th}(n,F)$ PROMPT NEUTRON MULTIPLICITY



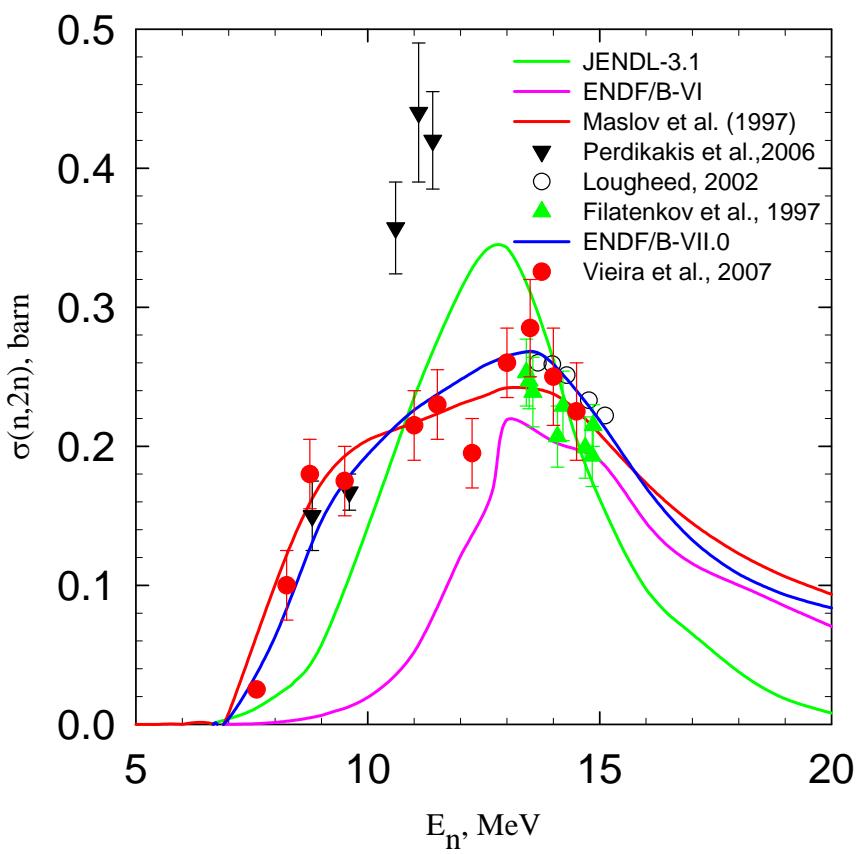
^{238}U FISSION CROSS SECTION



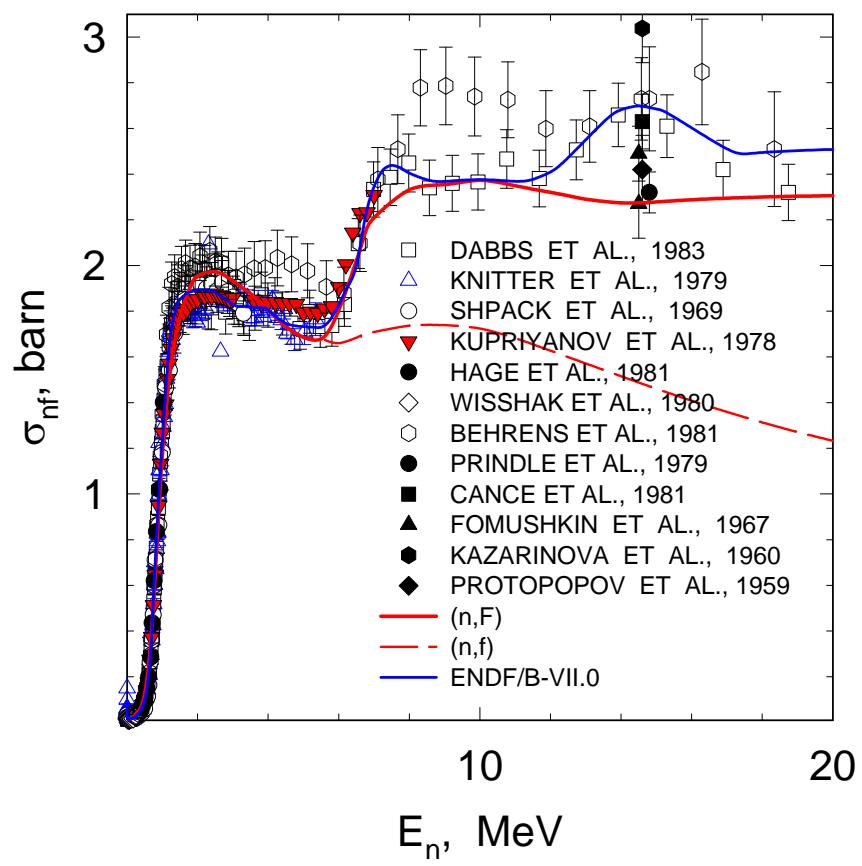
^{242}Cm FISSION CROSS SECTION



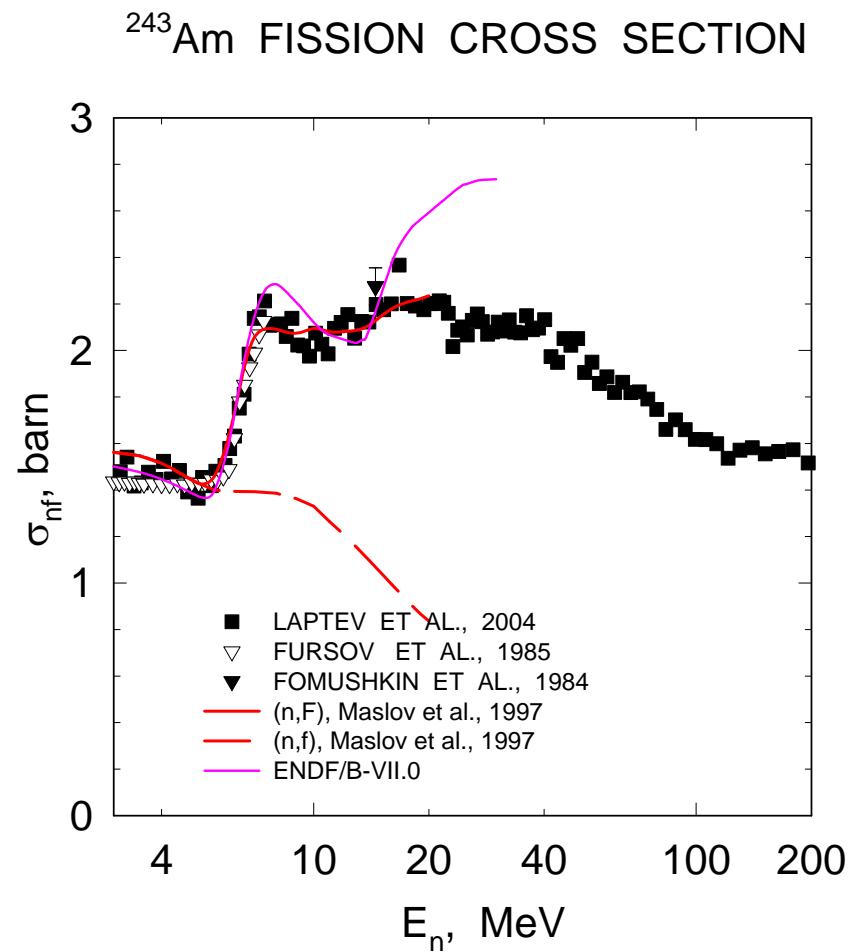
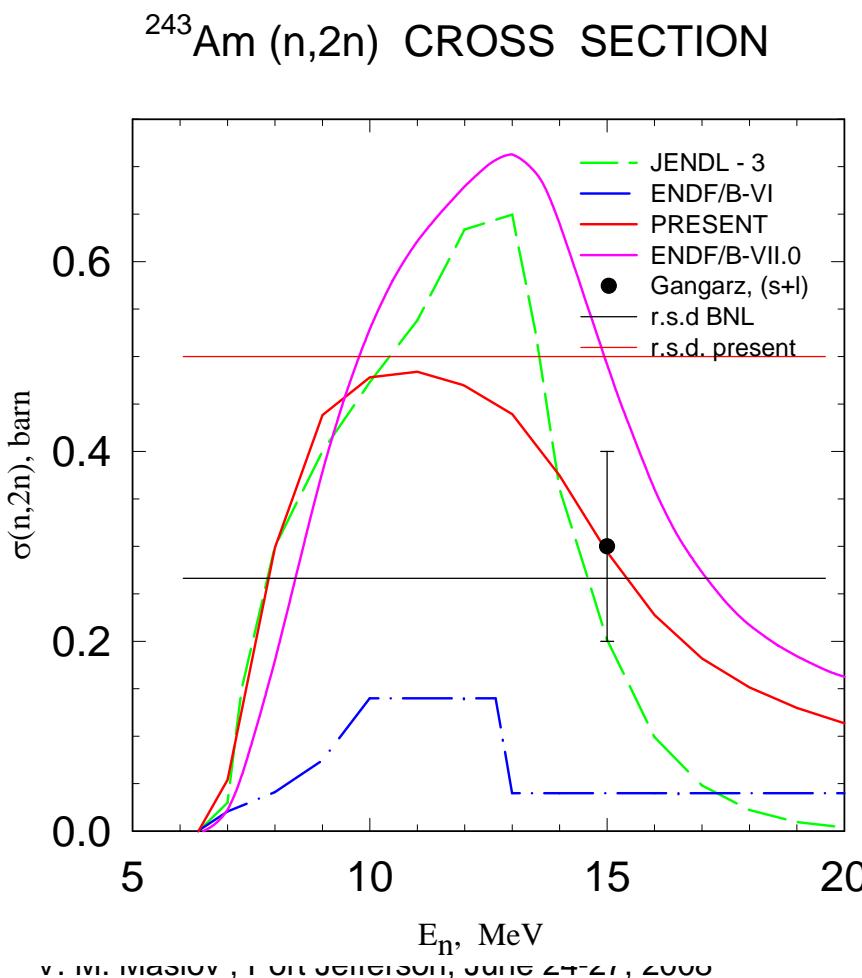
$^{241}\text{Am}(n,2n)$ CROSS SECTION



^{241}Am FISSION CROSS SECTION



$^{242}\text{gAm}/^{242}\text{mAm} \approx ^{236}\text{sNp}/^{236}\text{lNp} \approx 0.35$, Gancartz -
 $^{242}\text{gAm} = ^{242}\text{sAm}$

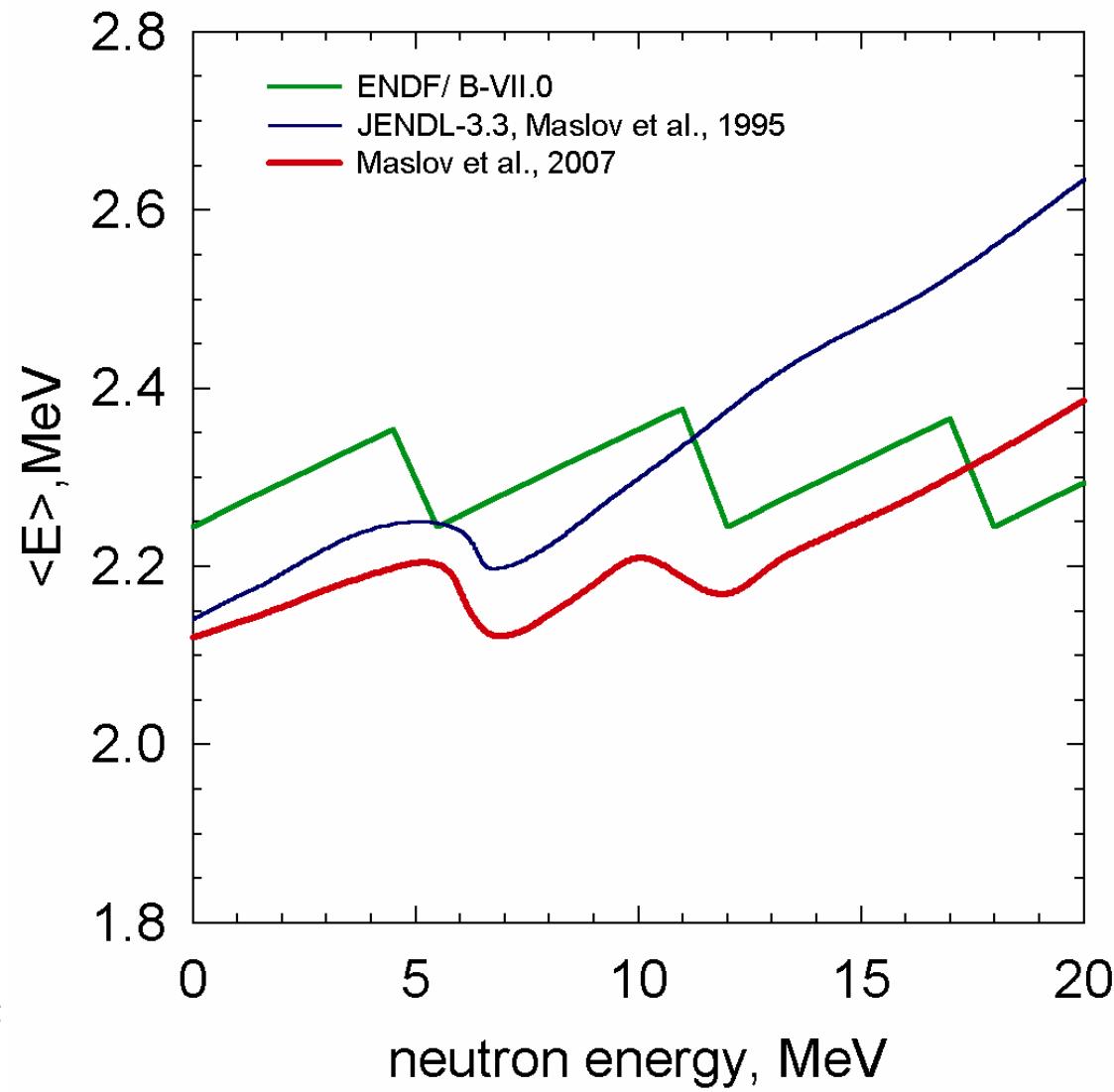


Prompt-fission neutron spectra

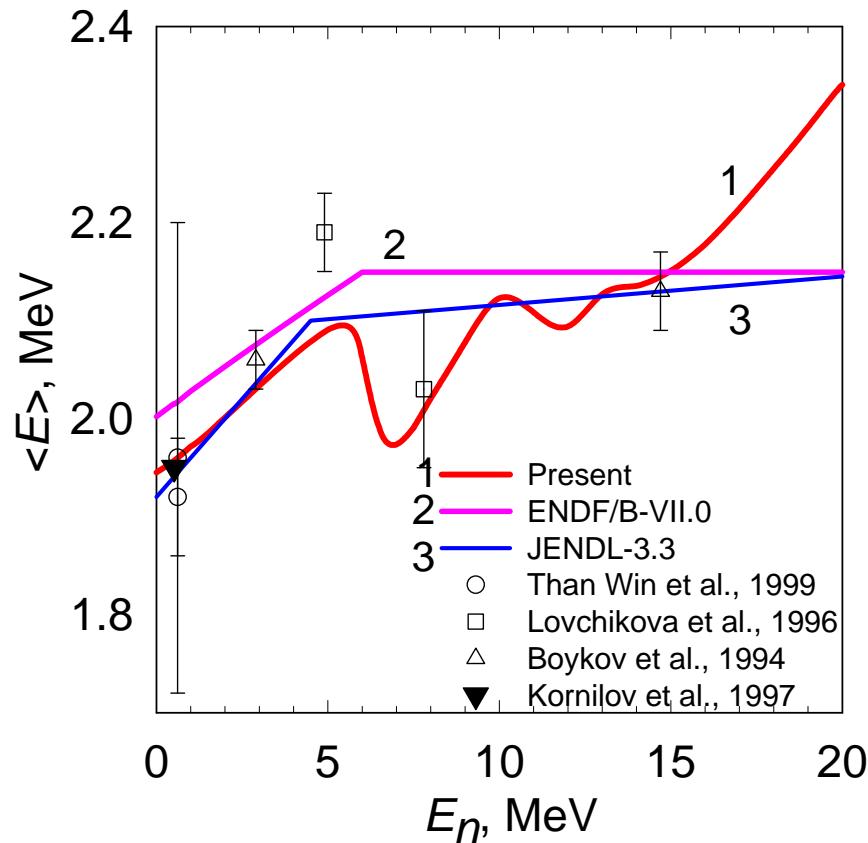
superposition of exclusive pre-fission (n,xnf) spectra and post-fission spectra

$$S_{A+2-x}(\varepsilon, E_n)$$

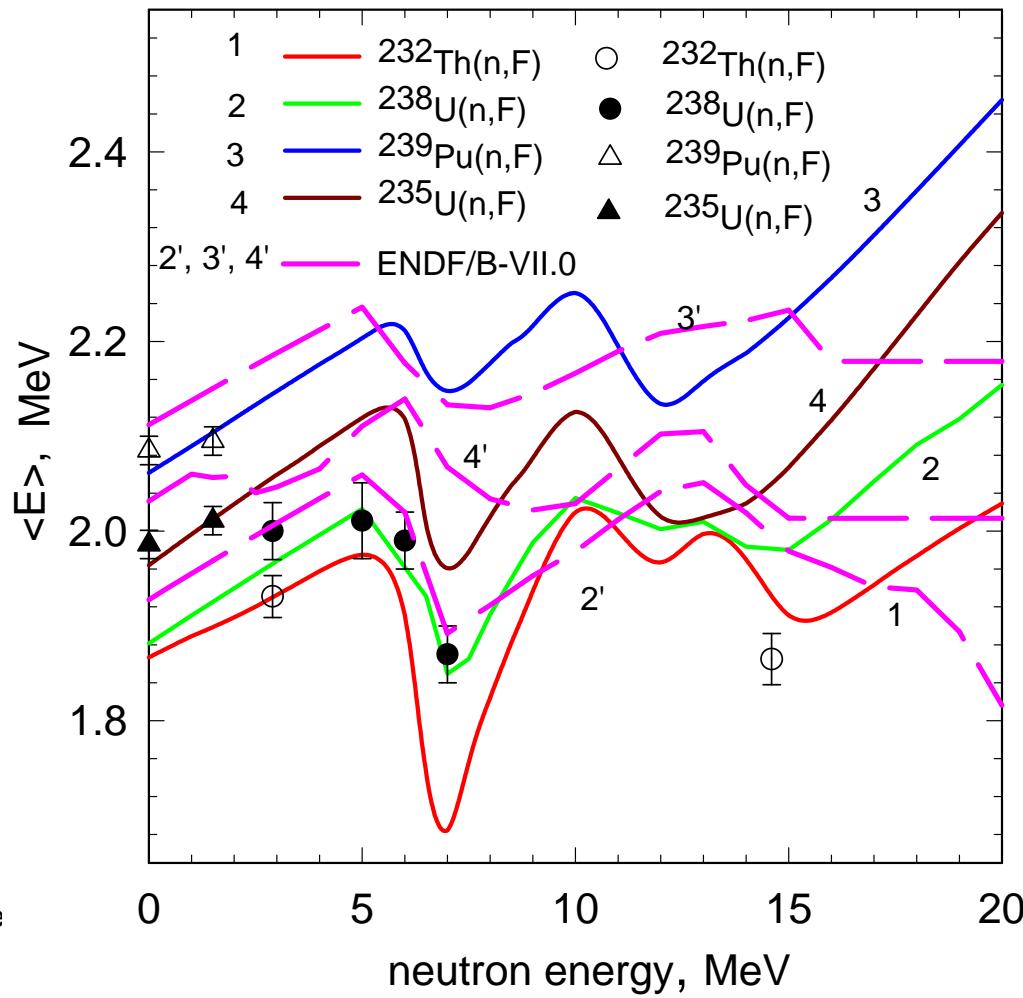
$$\begin{aligned} S(\varepsilon, E_n) = & \nu^{-1}(E_n) \{ \nu_1(E_n) \beta_1(E_n) S_{A+1}(\varepsilon, E_n) + \\ & \beta_2(E_n) [\nu_2(E_n) S_A(\varepsilon, E_n) + \frac{d\sigma_{\text{n2nnf}}^1(E_n)}{d\varepsilon}] + \\ & \beta_3(E_n) [\nu_3(E_n) S_{A-1}(\varepsilon, E_n) + (\frac{d\sigma_{\text{n2nnf}}^1(E_n)}{d\varepsilon} + \frac{d\sigma_{\text{n2nnf}}^2(E_n)}{d\varepsilon})] + \\ & \beta_4(E_n) [\nu_4(E_n) S_{A-2}(\varepsilon, E_n) + (\frac{d\sigma_{\text{n3nnf}}^1(E_n)}{d\varepsilon} + \frac{d\sigma_{\text{n3nnf}}^2(E_n)}{d\varepsilon} + \frac{d\sigma_{\text{n3nnf}}^3(E_n)}{d\varepsilon})] \} \end{aligned}$$

$^{241}\text{Am}(n, F), \langle E \rangle$ 

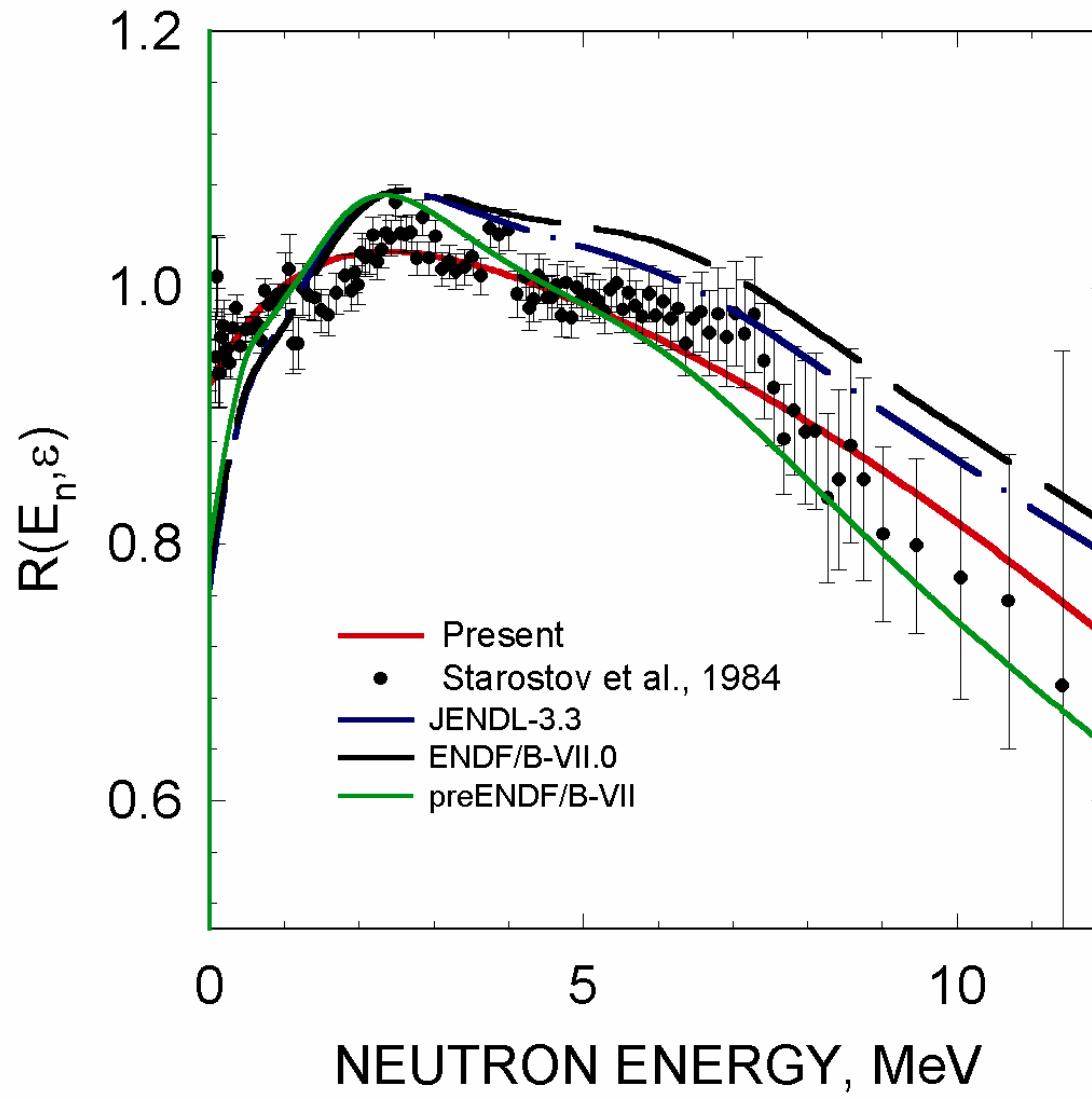
$^{237}\text{Np}(n,\text{F})$ AVERAGE ENERGY OF PFNS



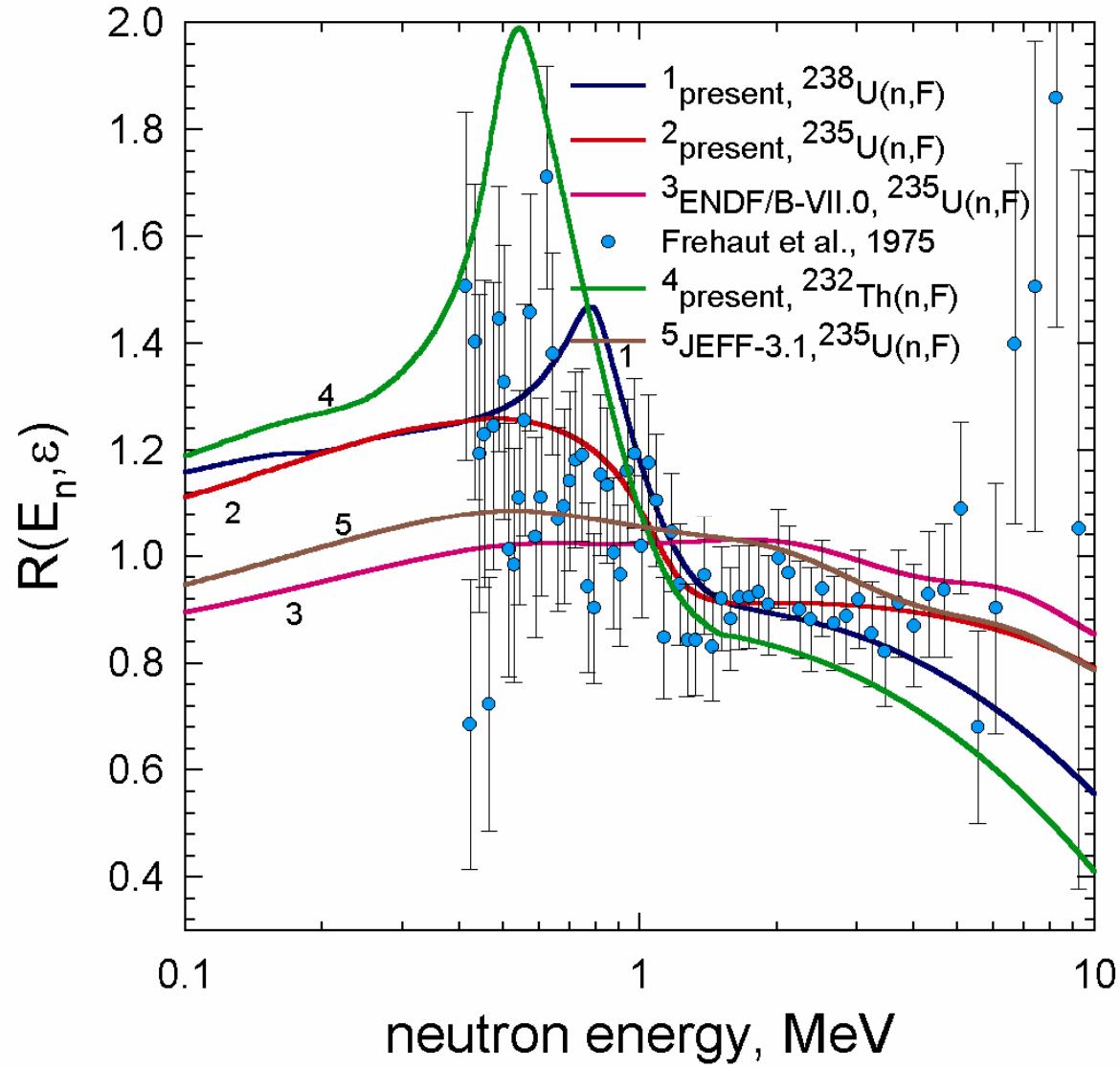
AVERAGE ENERGY OF PFNS



235_U(n, f), thermal



$^{235}\text{U}(n, f), E_n = 7 \text{ MeV}$



Conclusions

1. In a few examples of Pu, Am, Cm nuclides it is shown that **substituting model deficiency** uncertainties by enlarging the uncertainties of conventional nuclear model parameters **could** be avoided.
2. In a number of MA the **uncertainty estimation of CSS, PFNS** should be preceded with the robust neutron data re-evaluation. **Otherwise, in case of poorly investigated Np, Pu, Am, Cm targets the artificially large cross section uncertainty estimates will be unavoidable.**
3. Correlated evaluation of fission **CSS and PFNS for MA**, as for $^{237}\text{Np}(n, F)$ and $^{241}\text{Am}(n, F)$, supported by the PFNS measurement for $^{237}\text{Np}(n, F)$ by Taieb et al. [39].